

Partial Differential Equation

UNIT-I

Second Order Partial differential Equation

Origin of 2nd order P.D.E — linear differential eqn with constant co-efficients —
Methods of solving (linear) P.D.E — classification of 2nd order P.D.E — Canonical form —
Adjoint operator — Riemann method.

chapter: 2 Sec 2.1 to 2.5

UNIT-II

Elliptic Differential Equation

Occurrence of a Laplace and Poisson's Equation — Boundary value Problems —
separation of variables method — Laplace eqn in cylindrical and Spherical coordinates —
Dirichlet and Neumann problems for a circle and sphere

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UNIT-3

Parabolic Differential Equation and derivation

Occurrence of the diffusion eqn —
Boundary Condition — separation of variables method — Diffusion eqn in

cylindrical & spherical Co-ordinates

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Hyperbolic Differential Eqns

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One-dimensional wave eqn — Reduction to
Canonical form — D'Alembert's solution —
Separation of variables method — Periodic
- cylindrical - spherical
Co-ordinates — Duhamel's principle for
wave eqn.

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UNIT: 5

Integral Transform

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P.D.E — Diffusion eqn — wave equation —
Fourier Transforms — ~~their~~ application
to P.D.E — ^{diffusion eqn - wave eqn} Laplace equation

chapters: 6 Sec: 6.2 to 6.4

Text Book: -

P.D.E for Engineering → 2nd Edition.

J.N. Sharma & Kesh

UNIT-I

Second order P.D.E

✓ Origin of 2nd order partial differential eqn:-
 Consider the function $z = f(u) + g(v) + w$

Soln. where, f and g are arbitrary function of u and v respectively and u, v and w are functions of x and y .

we write,

$$p = \frac{\partial z}{\partial x} \quad ; \quad q = \frac{\partial z}{\partial y} \quad ; \quad r = \frac{\partial^2 z}{\partial x^2}$$

$$s = \frac{\partial^2 z}{\partial x \partial y} \quad , \quad t = \frac{\partial^2 z}{\partial y^2}$$

Diff. (1) w.r. to x and y respectively.

we get.

$$p = f'(u) u_x + g'(v) v_x + w_x \quad \text{--- (2)}$$

$$q = f'(u) u_y + g'(v) v_y + w_y \quad \text{--- (3)}$$

$$r = f''(u) (u_x)^2 + f'(u) u_{xx} + g''(v) (v_x)^2 + g'(v) v_{xx} + w_{xx} \quad \text{--- (4)}$$

$$s = f''(u) u_x u_y + g''(v) v_x v_y + f'(u) u_{xy} + g'(v) v_{xy} + w_{xy} \quad \text{--- (5)}$$

$$t = f''(u) u_y^2 + f'(u) u_{yy} + g''(v) v_y^2 + g'(v) v_{yy} + w_{yy} \quad \text{--- (6)}$$

Eliminating the arbitrary quantities f, g, f'' and g'' in (a) to (b)

we get

$$\begin{vmatrix} P - w_x & U_x & V_x & 0 & 0 \\ q - w_y & V_y & V_y & 0 & 0 \\ r - w_{xx} & U_{xx} & V_{xx} & U_x^2 & V_x^2 \\ s - w_{xy} & U_{xy} & V_{xy} & U_x V_y & V_x V_y \\ t - w_{yy} & U_{yy} & V_{yy} & U_y^2 & V_y^2 \end{vmatrix} = 0$$

If we expand in terms of first column, we get

$$\boxed{R_r + S_s + T_t + P_p + Q_q = W} \quad (7)$$

where,

R, S, T, P, Q & W are known functions of x and y .

\therefore Equation (1) is solution of eqn (7).

\therefore It is a P.D.E of second order.

PROBLEM:-

(1) If $U = f(x+iy) + g(x-iy)$ where f & g are arbitrary functions s.t.

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

Soln:-

Given,

$$U = f(x+iy) + g(x-iy) \quad \text{--- (1)}$$

Diff (1) p.w.r. to x & y respectively.

$$\frac{du}{dx} = f'(x+iy) + g'(x-iy)$$

$$\frac{d^2u}{dx^2} = f''(x+iy) + g''(x-iy)$$

$$\frac{du}{dy} = f'(x+iy)i + g'(x-iy)(-i)$$

$$\frac{d^2u}{dy^2} = -f''(x+iy)(-1) + g''(x-iy)(-1)$$

$$(-i)(-i) = +i^2 = -1$$

$$\frac{d^2u}{dx^2} + \frac{d^2u}{dy^2} = 0$$

Q2 If f and g be arbitrary functions of the respectively arguments show that $u = f(x - vt + i\alpha y) + g(x - vt - i\alpha y)$ is a solution of $\frac{d^2u}{dx^2} + \frac{d^2u}{dy^2} = \frac{1}{c^2} \frac{d^2u}{dt^2}$ provided $\alpha^2 = 1 - \frac{v^2}{c^2}$

Soln:-

Given,

$$u = f(x - vt + i\alpha y) + g(x - vt - i\alpha y) \quad \text{--- (1)}$$

Diff (1) p.w.r to x & y respectively.

$$\frac{du}{dx} = f'(x - vt + i\alpha y) + g'(x - vt - i\alpha y)$$

$$\frac{d^2u}{dx^2} = f''(x - vt + i\alpha y) + g''(x - vt - i\alpha y) \quad \text{--- (2)}$$

$$\frac{du}{dy} = f'(x - vt + i\alpha y)(i\alpha) + g'(x - vt - i\alpha y)(-i\alpha)$$

$$\frac{d^2 U}{dy^2} = f''(x - vt + i\alpha y) (i\alpha)^2 + g''(x - vt - i\alpha y) \alpha^2$$

$$\frac{d^2 U}{dy^2} = -f''(x - vt + i\alpha y) \alpha^2 - g''(x - vt - i\alpha y) \alpha^2 \quad \text{--- (2)}$$

Eqn (1) Diff. P.w.r.to t

$$\frac{dU}{dt} = f'(x - vt + i\alpha y)(-v) + g'(x - vt - i\alpha y)v$$

$$\frac{d^2 U}{dt^2} = f''(x - vt + i\alpha y)v^2 + g''(x - vt - i\alpha y)v^2 \quad \text{--- (3)}$$

$$\text{(2) + (3)} \Rightarrow$$

$$\frac{d^2 U}{dx^2} + \frac{d^2 U}{dy^2} = f''(x - vt + i\alpha y) \alpha^2 + g''(x - vt - i\alpha y) \alpha^2$$

$$= (1 - \alpha^2) [f''(x - vt + i\alpha y) + g''(x - vt - i\alpha y)]$$

$$= (1 - \alpha^2) \frac{1}{v^2} \frac{d^2 U}{dt^2}$$

$$= \frac{1}{c^2} \frac{d^2 U}{dt^2}$$

$$\frac{d^2 U}{dx^2} + \frac{d^2 U}{dy^2} = \frac{1}{c^2} \frac{d^2 U}{dt^2}$$

Linear Partial differential equation with

Constant Co-efficient

An eqn is of the form

$$F(D, D')z = f(x, y) \quad \text{--- (1)}$$

where, $F(D, D')$ is a differential operator of a type

$$F(D, D') = \sum_r \sum_s C_{rs} D^r D'^s$$

$$2D^3 + 3D^2 + 5D$$

where,

$$D = \frac{\partial}{\partial x}$$

$$\& D' = \frac{\partial}{\partial y}$$

& C_{rs} are constant

which is called a Linear P.D.E with Constant Co-efficient.

Note:-

i) A solution of $F(D, D')z = 0$ is a solution called a Complementary function and $F(D, D')z = f(x, y)$ is called a Particular Integral.

∴ The general solution of (1) is

$$z = C \cdot F + P \cdot I$$

NOTE:-

We classify the operator $F(D, D') = 0$ into two types.

1) Reducible

2) Irreducible

✓ Reducible:-

⊗ The operator $F(D, D')$ is said to be reducible if it can be factorized into the linear factor of the type $D + aD' + b$ where a and b are constants.

Ex:-

$$1) F(D, D') = D^2 - D'^2$$

$$= (D + D')(D - D')$$

$$= (D + D' + 0)(D - D' + 0)$$

$$2) F(D, D') = D^2 + D$$

$$= D(D + 1)$$

$$= (D + 0D' + 0)(D + 0D' + 1)$$

⊗ Irreducible:-

2/17 The operator $F(D, D')$ is said to be irreducible if it is not reducible.

Ex:-

$$F(D, D') = D^2 + D$$

Reducible Equation:-

A P.D.E $F(D, D')z = f(x, y)$, is said to be a reducible, if $F(D, D')$ can be written as a product of linear factor in D & D' .

$$\text{i.e. } F(D, D') = \prod_{r=1}^n (\alpha_r D + \beta_r D' + \gamma_r)$$

If the equation is not reducible then it is called irreducible equation.

Thm: 2

⊙ If $\alpha_r D + \beta_r D' + \gamma_r$ is a factor of $F(D, D')$ & $\phi_r(\Sigma)$ is an arbitrary function of the single variable Σ then

$$U_r = \exp\left(\frac{-\gamma_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y), \alpha_r \neq 0 \text{ is}$$

the solution of the equation $F(D, D')z = 0$

Proof:

The Given Equation is

$$F(D, D')z = 0 \quad \text{--- (1)}$$

we have to prove

$$U_T = \exp\left(-\frac{\gamma_T x}{\alpha_T}\right) \phi_T'(\beta_T x - \alpha_T y) \text{ is a}$$

Solution of equation (1)

It is enough to prove that

$$\Gamma(D, D')U_T = 0$$

Claim: - $(\alpha_T D + \beta_T D' + \gamma_T)U_T = 0$

Let $\alpha_T D(U_T) = \alpha_T \frac{d}{dx}(U_T)$

$$= \alpha_T \frac{d}{dx} \left(\exp\left(-\frac{\gamma_T x}{\alpha_T}\right) \phi_T'(\beta_T x - \alpha_T y) \right)$$

$$= \alpha_T \left[\exp\left(-\frac{\gamma_T x}{\alpha_T}\right) \phi_T'(\beta_T x - \alpha_T y) \beta_T \right.$$

$$\left. + \phi_T'(\beta_T x - \alpha_T y) \exp\left(-\frac{\gamma_T x}{\alpha_T}\right) \left(-\frac{\gamma_T}{\alpha_T}\right) \right]$$

$$\alpha_T D(U_T) = \exp\left(-\frac{\gamma_T x}{\alpha_T}\right) \phi_T'(\beta_T x - \alpha_T y) \beta_T \alpha_T$$

$$- \phi_T'(\beta_T x - \alpha_T y) \exp\left(-\frac{\gamma_T x}{\alpha_T}\right) \gamma_T \quad (2)$$

$$\beta_T D'(U_T) = \beta_T \frac{d}{dy}(U_T)$$

$$= \beta_T \frac{d}{dy} \left(\exp\left(-\frac{\gamma_T x}{\alpha_T}\right) \phi_T'(\beta_T x - \alpha_T y) \right)$$

$$= \beta_T \exp\left(-\frac{\gamma_T x}{\alpha_T}\right) \phi_T'(\beta_T x - \alpha_T y) (-\alpha_T)$$

$$\begin{aligned} \alpha_r D U_r + \beta_r D' U_r &= \phi_r (\beta_r x - \alpha_r y) \exp\left(-\frac{\gamma_r x}{\alpha_r}\right) (-\gamma_r) \\ &= -\phi_r (\beta_r x - \alpha_r y) \exp\left(-\frac{\gamma_r x}{\alpha_r}\right) \gamma_r \\ &= -U_r \gamma_r \end{aligned}$$

$$\alpha_r D U_r + \beta_r D' U_r + \gamma_r U_r = 0$$

$$(\alpha_r D + \beta_r D' + \gamma_r) U_r = 0$$

Since $\alpha_r D + \beta_r D' + \gamma_r$ is a factor of $F(D, D')$

we have write

$$F(D, D') = F_1(D, D') (\alpha_r D + \beta_r D' + \gamma_r)$$

where $F_1(D, D')$ are some other factors of $F(D, D')$,

$$\begin{aligned} \therefore F(D, D') U_r &= F_1(D, D') (\alpha_r D + \beta_r D' + \gamma_r) U_r \\ &= F_1(D, D') (0) = 0 \end{aligned}$$

$$F(D, D') U_r = 0$$

$$\therefore U_r = \exp\left(-\frac{\gamma_r x}{\alpha_r}\right) \phi_r (\beta_r x - \frac{\alpha_r y}{\alpha_r}) \text{ is a}$$

Solution of $F(D, D') z = 0$..

Method of solving linear P.D.E

The solution of reducible equation

$$F(D, D')z = f(x, y) \quad \text{--- (1)}$$

Be a given P.D.E

$$\text{If } f(D, D')z = \sum_{r=1}^n [\alpha_r D + \beta_r D' + \gamma_r]z$$

If z satisfies

$$(\alpha_r D + \beta_r D' + \gamma_r)z = 0 \quad r = 0, 1, 2, \dots, n$$

$$\alpha_r \frac{dz}{dx} + \beta_r \frac{dz}{dy} + \gamma_r z = 0$$

$$\alpha_r P + \beta_r Q = -\gamma_r z$$

which is a Lagrangian type of a P.D.E

$$\text{i.e.) } Pp + Qq = R$$

$$\frac{dx}{P} = \frac{dy}{Q} = \frac{dz}{R}$$

$$\frac{dx}{\alpha_r} = \frac{dy}{\beta_r} = \frac{dz}{-\gamma_r z} \quad \text{--- (2)}$$

$$\text{Take } \frac{dx}{\alpha_r} = \frac{dy}{\beta_r} \Rightarrow \beta_r dx = \alpha_r dy$$

$$\int \beta_r dx = \int \alpha_r dy + C_r$$

$$\alpha \beta_r - y \alpha_r = C_r \quad \text{--- (3)}$$

Take

$$\frac{dx}{\alpha_r} = \frac{dz}{-\gamma_r z}$$

$$-\frac{\gamma_r}{\alpha_r} dz = \frac{dz}{z}$$

Integ $-\frac{\gamma_r}{\alpha_r} x = \log z - \log A_r$ [where, A_r Constant]

$$-\frac{\gamma_r x}{\alpha_r} = \log \left(\frac{z}{A_r} \right)$$

$$\exp \left(-\frac{\gamma_r x}{\alpha_r} \right) = \frac{z}{A_r}$$

$$z = A_r \exp \left(-\frac{\gamma_r x}{\alpha_r} \right)$$

$$z = \phi_r (C_r) \exp \left(-\frac{\gamma_r x}{\alpha_r} \right)$$

$$z = \phi_r (x \beta_r - y \alpha_r) \exp \left(-\frac{x \gamma_r}{\alpha_r} \right)$$

we assume $\alpha_r \neq 0$

\therefore The Complementary function $\left. \begin{array}{l} \\ \end{array} \right\} = y = \sum_{r=1}^n \phi_r (x \beta_r - y \alpha_r) \exp \left(\frac{x \gamma_r}{\alpha_r} \right)$

Particular case:-

If $\alpha_r = 0$ then

$$\textcircled{2} \Rightarrow \frac{dx}{0} = \frac{dy}{\beta_r} = \frac{dz}{-\gamma_r z}$$

1st & 2nd term

$$\frac{dx}{0} = \frac{dy}{\beta_r}$$

$$\beta_r dx = 0$$

$$\int \beta_r dx = 0$$

$$\beta_r x = C_r$$

nd & 3rd term

$$\frac{dy}{\beta_r} = \frac{dz}{-\gamma_r z}$$

$$\int \frac{-\gamma_r}{\beta_r} dy = \int \frac{dz}{z}$$

$$-\frac{\gamma_r}{\beta_r} y = \log z - \log A_r$$

$$\exp\left(-\frac{\gamma_r}{\beta_r} y\right) = \frac{z}{A_r}$$

$$z = A_r \exp\left(\frac{-\gamma_r y}{\beta_r}\right)$$

$$z = \phi_r(C_r) \exp\left(\frac{-\gamma_r y}{\beta_r}\right)$$

$$z = \phi_r(\beta_r x) \exp\left(\frac{-\gamma_r y}{\beta_r}\right)$$

The above two cases are applicable.

When there is no repeated factors of the type $\alpha_r D + \beta_r D' + \gamma_r$:

Repeated Factors:-

Let P.D.E. $F(D, D')z = f(x, y)$ have

repeated factors.

i.e) suppose $(\alpha_r D + \beta_r D' + \gamma_r)^2$ is a factor of $F(D, D')$.

Then, we have

$$(\alpha_r D + \beta_r D' + \gamma_r)^2 z = 0$$

$$\text{Let } (\alpha_r D + \beta_r D' + \gamma_r) z = z_1 \quad \text{--- (A)}$$

where z_1 is a function of x & y .
such that

$$(\alpha_r D + \beta_r D' + \gamma_r) z_1 = 0$$

W.K.T

The above eqn has a solution of the form

$$z_1 = \exp\left(\frac{-\gamma_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y), \quad \alpha_r \neq 0$$

sub z_1 in equ (A)

$$(\alpha_r D + \beta_r D' + \gamma_r) z = \exp\left(\frac{-\gamma_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y)$$

$$\alpha_r D z + \beta_r D' z + \gamma_r z = \exp\left(\frac{-\gamma_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y)$$

$$\alpha_r p + \beta_r q + \gamma_r z = \exp\left(\frac{-\gamma_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y)$$

$$\alpha_r p + \beta_r q = \exp\left(\frac{-\gamma_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y) - \gamma_r z$$

$$\frac{dx}{\alpha_r} = \frac{dy}{\beta_r} = \frac{dz}{\exp\left(\frac{-\gamma_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y) - \gamma_r z}$$

Take (1) & (2)

$$\frac{dx}{\alpha_r} = \frac{dy}{\beta_r}$$

$$\text{Jing} \Rightarrow \beta_r x = \alpha_r y + C_r$$

$$C_r = \beta_r x - \alpha_r y$$

Also ① & ③

$$\frac{dx}{\alpha_r} = \frac{dz}{\exp\left(-\frac{\beta_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y) - \beta_r z}$$

$$\frac{1}{\alpha_r} \exp\left(-\frac{\beta_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y) - \frac{\beta_r z}{\alpha_r} = \frac{dz}{dx}$$

$$\frac{dz}{dx} + \frac{\beta_r z}{\alpha_r} = \frac{1}{\alpha_r} \exp\left(-\frac{\beta_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y)$$

This is of the form of L.D.E

$$\frac{dy}{dx} + py = Q$$

$$\text{Here, } y = z, p = \frac{\beta_r}{\alpha_r} \text{ \& } Q = \frac{1}{\alpha_r} \exp\left(-\frac{\beta_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y)$$

$$\Rightarrow \text{Integrating factor, I.F} = e^{\int p dx}$$

$$= e^{\int \frac{\beta_r}{\alpha_r} dx}$$

$$= \exp\left(\frac{\beta_r x}{\alpha_r}\right)$$

The solution is

$$y e^{\int p dx} = \int Q (e^{\int p dx}) dx + c$$

$$z \exp\left(\frac{\beta_r x}{\alpha_r}\right) = \int \left[\frac{1}{\alpha_r} \exp\left(-\frac{\beta_r x}{\alpha_r}\right) \phi_r(\beta_r x - \alpha_r y) \exp\left(\frac{\beta_r x}{\alpha_r}\right) \right] dx + c$$

$$= \int \frac{1}{\alpha_r} \phi_r (\beta_r x - \alpha_r y) dx + C$$

where $C_r = \beta_r x - \alpha_r y$

$$= \frac{\phi_r C_r^{\alpha_r}}{\alpha_r} + \phi_r(C_r)$$

is a const.

$$\therefore Z = \exp\left(-\frac{\gamma_r x}{\alpha_r}\right) \left(\frac{\phi_r C_r^{\alpha_r}}{\alpha_r} + \phi_r(C_r) \right)$$

The procedure can be generalized upto any order of repetition factor, adding this to the some of other solution corresponding to the linear factors without repetition.

\therefore we get the required complementary function.

Solution of irreducible equation with

constant co-efficient:-

Let $F(D, D')z = f(x, y)$ — (1) be a irreducible linear P.D.E with constant

co-efficient.

$$\text{Let } F(D, D') = F_1(D, D') F_2(D, D') \text{ — (2)}$$

where F_1 is irreducible and F_2 is reducible.

$$\textcircled{1} \Rightarrow F_1(D, D') F_2(D, D') z = f(x, y)$$

Then, the solution of corresponding linear factor of $F_2(D, D')$ can be obtained as

$$z = \exp\left(-\frac{P_1 x}{d_1}\right) \phi_1(P_1 x + \bar{d}_1 y), \quad d_1 \neq 0$$

(or)

$$z = \exp\left(-\frac{P_1 x}{d_1}\right) \phi_1(P_1 x) \quad \text{if } d_1 = 0.$$

now, to find the soln of irreducible

$F_1(D, D')$.

suppose, $z = e^{ax+by}$ is a soln of

$$F_1(D, D') z = 0 \quad \textcircled{3}$$

$$\therefore F_1(D, D') e^{ax+by} = F_1(a, b) e^{ax+by}$$

the above equation is must be vanish & we impose the condition that $F_1(a, b) = 0$

$$z = \sum_r C_r e^{a_r x + b_r y} \quad \text{is a Complementary}$$

function of equation $\textcircled{3}$

$$F_1(a_r, b_r) = 0, \quad r = 1, 2, \dots$$

Rules for finding Complementary function:-

Consider the equation

$$\frac{d^2 z}{dx^2} + a_1 \frac{d^2 z}{dx dy} + a_2 \frac{d^2 z}{dy^2} = 0 \text{ it can be}$$

written as $(D^2 + a_1 D D' + a_2 D'^2) z = 0$ — (1)

$$\text{where } D = \frac{\partial}{\partial x}, \quad D' = \frac{\partial}{\partial y}$$

Auxiliary equation:-

$$\text{(1)} \Rightarrow \frac{D^2}{D'^2} + \frac{a_1 D}{D'} + a_2 = 0$$

$$m^2 + a_1 m + a_2 = 0 \text{ — (2)}$$

where $m = D/D'$

which is called auxiliary equation

Let m_1 & m_2 be the roots of (2).

Case (i):-

$$\text{when } m_1 \neq m_2 \quad \begin{cases} m - m_1 = 0 \\ m - m_2 = 0 \end{cases}$$

$$\begin{cases} m = m_2 = m_1 \\ \frac{\partial}{\partial x} = m_1 \end{cases}$$

$$\text{(1)} \Rightarrow (D - m_1 D') (D - m_2 D') z = 0 \text{ — (3)}$$

Now, the solution of $(D - m_2 D') z = 0$

is also a solution of equation (3)

$$(D - m_2 D')z = 0$$

$$Dz - m_2 D'z = 0$$

$$p - m_2 q = 0$$

which is Lagrangian form

$$\therefore \frac{dx}{1} = \frac{dy}{-m_2} = \frac{dz}{0}$$

1st & 2nd term,

$$\frac{dx}{1} = \frac{dy}{-m_2}$$

$$-m_2 x = y + C_1$$

$$C_1 = y + m_2 x$$

Integ

$$0 = \int dz$$

$$z = C_2 = f_2(C_1)$$

$\therefore z = f_2(m_2 x + y)$ is a solution of

$(D - m_2 D')z = 0$, where f_2 is arbitrary

function.

||
||
||

A solution of $(D - m_1 D')z = 0$ is also

solution of ③ proceeding the same way.

$$z = f_1(y + m_1 x)$$

where f_1 is another arbitrary function.

∴ The C.F of equation (1) is given

by

$$z = f_1(y + m_1 x) + f_2(y + m_2 x) \dots$$

Case (ii): -

$$m_1 = m_2 = m$$

$$(1) \Rightarrow (D - mD')z = 0$$

$$(D - mD')(D - mD')z = 0$$

$$\text{Let } u = (D - mD')z$$

$$(D - mD')u = 0$$

The solution $u = f(y + mx)$ by case (i)

but

$$u = (D - mD')z$$

$$\therefore f(y + mx) = (D - mD')z$$

$$= p - mq$$

by Lagrange's form

$$\frac{dx}{1} = \frac{dy}{-m} = \frac{dz}{f(y + mx)}$$

(1) & (2) term

$$\frac{dx}{1} = \frac{-dy}{m}$$

Integ

$$x = -\frac{1}{m}y + c$$

$$c_1 = x + \frac{y}{m} = mx + y$$

(1) & 3rd term

$$\frac{dx}{1} = \frac{dz}{f(y+mx)} \Rightarrow \frac{dx}{1} = \frac{dz}{f(c_1)}$$

Integ $f(c_1)x + c_2 = z$

$$z = f(y+mx)x + c_2$$

∴ The C.F of equ (1) is

Generalized the result of case (i) and case (ii)
$$z = F_1(y+mx) + x f_2(y+mx)$$

If the ~~roots~~ ^{rules} of aux^y auxiliary equation are m_1, m_2, \dots (distinct)

Then the Complementary function

$$C.F = f_1(y+m_1x) + f_2(y+m_2x) + \dots$$

If the two roots of auxiliary eqn is equal $m_1 = m_2 = m$.

$$\text{Then, } C.F = f_1(y+mx) + x f_2(y+mx)$$

If three roots of auxiliary equation are equal $m_1 = m_2 = m_3 = m$.

$$C.F = f_1(y+mx) + x f_2(y+mx) + x^2 f_3(y+mx)$$

Rules for finding particular integral:-

Consider the equation $\phi(D, D')z = F(x, y)$

where $\phi(D, D') = D^n + a_1 D^{n-1} + a_2 D^{n-2} + \dots$

Then

$$P.I = \frac{1}{\phi(D, D')} F(x, y)$$

Case i):-

i) when $F(x, y) = e^{ax+by}$

$$\text{Then, } P.I = \frac{1}{\phi(a, b)} e^{ax+by}, \quad \phi(a, b) \neq 0$$

If $\phi(a, b) = 0$, then this case is failure.

ii) ~~is~~ when $F(x, y) = \sin(ax+by)$

$$P.I = \frac{1}{\phi(\omega^2, D\omega, D'^2)} \sin(ax+by)$$

$$= \frac{1}{\phi(-a^2, -ab, -b^2)} \sin(ax+by), \quad \phi(-a^2, -ab, -b^2) \neq 0$$

If $\phi(-a^2, -ab, -b^2) = 0$ then this case is failure.

A similar rule for $F(x, y) = \cos(ax+by)$.

iii) $F(x, y) = x^m y^n$, where m & n are +ve integers.

Then P.I = $\frac{1}{\phi(D, D')} x^m y^n$

= $(\phi(D, D'))^{-1} x^m y^n$

If $m < n$ we expand as binomially $[\phi(D, D')]^{-1}$ in power of $\frac{D}{D'}$ and $m > n$ we expand $[\phi(D, D')]^{-1}$ in power of D'/D

iv) Also $\frac{1}{D} F(x, y) = \int F(x, y) dx$, $y = \text{constant}$

$\frac{1}{D'} F(x, y) = \int F(x, y) dy$, $x = \text{constant}$

Case ii): General method for to find P.I

The General method is applicable to all cases where $F(x, y)$ is not of form if on above. Now, $F(D, D')$ can be factorized in general into linear factors.

P.I = $\frac{1}{\phi(D, D')} F(x, y)$

= $\frac{1}{(D - m_1 D') (D - m_2 D') \dots (D - m_n D')} F(x, y)$

Let us evaluate $\frac{1}{(D - m D')} F(x, y) = Z$

Consider, $F(x, y) = Z (D - m D')$

$$\frac{dx}{-} = \frac{dy}{-m} = \frac{dz}{F(x,y)}$$

① & ②

$$\frac{dx}{-} = \frac{dy}{-m}$$

Integ $-m \int \frac{dx}{-} = \int dy$
 $C_1 = y + mx$

① & ③rd term

$$dz = \frac{dz}{F(x,y)}$$

$$F(x,y) dz = dz$$

$$\int F(x, C_1 - mx) dz = \int dz$$

$$\frac{1}{(D - mD')} F(x,y) z = \int F(x, C_1 - mx) dx$$

where C_1 is replaced by $y + mx$

after integration.

Solve the equation $\frac{d^3 z}{dx^3} - 2 \frac{d^3 z}{dx^2 dy} - \frac{d^3 z}{x dy^2} +$

$2 \frac{d^3 z}{dy^3} = e^{x+y}$

Soln:-

The given eqn is written as

$$(D^3 - 2D^2 D' - D D'^2 + 2D'^3) z = e^{x+y} \quad \text{--- (1)}$$

Here,

$$F(D, D') = D^3 - 2DD'^2 - DD'^2 + 2D'^3$$

$$= D^2(D - 2D') - (D'^2[D - 2D'])$$

$$= (D^2 - D'^2)(D - 2D')$$

$$F(D, D') = (D + D')(D - D')(D - 2D') \quad \text{--- (2)}$$

Now,

$$(D - 2D')z = 0$$

$$p - 2q = 0$$

In Lagrange's form

$$\frac{dx}{1} = \frac{dy}{-2} = \frac{dz}{0}$$

1st & 2nd term

$$\frac{dx}{1} = \frac{dy}{-2}$$

$$\int -2 dx = \int dy$$

$$-2x + c = y$$

$$c = y + 2x$$

1st and 3rd

$$\frac{dx}{1} = \frac{dz}{0}$$

$$\int dx = \int dz$$

$$z = \phi(cc) = \phi_1(y + 2x) \text{ is a soln of (2)}$$

Now,

$$(D + D')z = 0$$

$$P + Q = 0$$

In Lagrange's form

$$\frac{dx}{1} = \frac{dy}{1} = \frac{dz}{0}$$

1st & 2nd term

$$\frac{dx}{1} = \frac{dy}{1}$$

$$\int dx = \int dy$$

$$x + c = y$$

$$c = y - x$$

1st & 3rd term

$$\frac{dx}{1} = \frac{dz}{0}$$

$$\int dx = \int dz$$

$$z = \phi_2(c)$$

$$z = \phi_2(y - x)$$

$$C.F = \phi_1(y + x) + \phi_2(y - x) + \phi_3(x + y)$$

$$(P + D')z = 0$$

$$P - Q = 0$$

In Lagrange's form

$$\frac{dx}{1} = \frac{dy}{-1} = \frac{dz}{0}$$

1st & 2nd term

$$\frac{dx}{1} = \frac{dy}{-1}$$

$$\int -dx = \int dy$$

$$-x + c = y$$

$$c = y + x$$

1st & 3rd term

$$\frac{dx}{1} = \frac{dz}{0}$$

$$\int dx = \int dz$$

$$\phi_3(c) = z$$

$$z = \phi_3(y + x)$$

Particular Integral :-

$$P \cdot I = \frac{1}{\phi(D, D')} F(x, y)$$

$$= \frac{1}{(D+D')(D-D')(D-D)} e^{x+y}$$

$$= \frac{1}{2(D-D')(1-D)} e^{x+y}$$

$$P \cdot I = -\frac{1}{2(D-D')} e^{x+y} \quad \text{--- (3)}$$

Next,

$$w = \frac{1}{D-D'} e^{x+y}$$

$$(D-D')w = e^{x+y}$$

$$p - q = e^{x+y}$$

by Lagrange's form

$$\frac{dx}{1} = \frac{dy}{-1} = \frac{dw}{e^{x+y}}$$

(1) & (2)

$$\frac{dx}{1} = \frac{dy}{-1}$$

Integ

$$\int -dx = \int dy$$

$$c_1 = x+y$$

(1) & (3)

$$\frac{dx}{1} = \frac{dw}{e^{x+y}}$$

$$e^{x+y} \cdot dx = dw$$

$$\int e^{x+y} \cdot dx = \int dw$$

$$e^{x+y} \cdot x = w$$

$$w = e^{x+y} \cdot x$$

(3) \Rightarrow

$$P.I = -\frac{1}{2} w$$

$$P.I = -\frac{1}{2} \cdot e^{x+y} \cdot x$$

\therefore The general soln is

$$z = C.F + P.I$$

$$z = \phi_1(y+2x) + \phi_2(y-x) + \phi_3(x+y) - \frac{1}{2} x \cdot e^{x+y}$$

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solve $\frac{d^4 z}{dx^4} + \frac{d^4 z}{dy^4} = 2 \frac{d^4 z}{dx^2 dy^2}$

Soln:-

$$\text{Given } D^4 + D'^4 = 2D^2 D'^2$$

$$D^4 + D'^4 - 2D^2 D'^2 = 0$$

$$(D^2 - D'^2)^2 = 0$$

$$((D+D')(D-D'))^2 = 0$$

$$(D+D')^2 (D-D')^2 = 0$$

$$(D+D')^2 (D+D') (D-D') (D-D') = 0$$

$$i) (D+D')^2 z = 0 \quad \text{--- (2)}$$

$$\text{Let } (D+D') z = z_1 \quad \text{--- (3)}$$

$$\therefore (D+D') z_1 = 0$$

$$P+Q = 0$$

by Lagrange's form

$$\frac{dx}{1} = \frac{dy}{1} = \frac{dz_1}{0}$$

(1) & (2)

$$\int dx = \int dy$$

$$x - y = c_1$$

(1) & (3) $\frac{dx}{1} = \frac{dz_1}{0}$

$$z_1 = c = \phi_1(x)$$

$$z_1 = \phi_1(x - y)$$

$$(3) \Rightarrow (D+D') z = z_1 = \phi_1(x - y)$$

$$P+Q = \phi_1(x - y)$$

by Lagrange's form,

$$\frac{dx}{1} = \frac{dy}{1} = \frac{dz}{\phi_1(x - y)}$$

(1) & (2)

$$\int \frac{dx}{1} = \int \frac{dy}{1}$$

$$x - y = c_2$$

(1) & (3)

$$\int \frac{dx}{1} = \int \frac{dz}{\phi_1(x-y)}$$

$$\int \phi_1(c_2) dx = \int dz$$

$$x \phi_1(x-y) + \phi_2(c_2) = z.$$

$$\therefore z = x \phi_1(x-y) + \phi_2(x-y).$$

ii) $(D-D')^2 = 0$

Let $(D-D')z = z_2$

$$\therefore (D-D')z_2 = 0$$

by Lagrange's form $P=Q=0$

$$\frac{dx}{1} = \frac{dy}{-1} = \frac{dz_2}{0}$$

(1) & (2)

$$\frac{dx}{1} = \frac{dy}{-1}$$

$$\int -dx = \int dy$$

$$-x + c = y$$

$$c_3 = x+y.$$

(1) & (3)

$$\frac{dx}{1} = \frac{dz_2}{0}$$

$$\int dx = \int dz_2$$

$$z_2 = \phi_3(c_3)$$

$$z_2 = \phi_3(x+y)$$

~~ii) $(D-D')$ $\neq 0$~~
 ~~$(D-D')$ $\neq 0$~~

$$(A) \Rightarrow (D-D')z = z_2 = \phi_3(x+y)$$

$$(D-D')z = \phi_3(x+y)$$

$$P-Q = \phi_3(x+y)$$

By Lagrange form

$$\frac{dx}{1} = \frac{dy}{-1} = \frac{dz_2}{\phi_3(x+y)}$$

① a ②

$$\frac{dx}{1} = \frac{dy}{-1}$$

$$\int -dx = \int dy$$

$$-x + c_A = y$$

$$c_A = x+y$$

1st \int 2nd

$$\frac{dx}{1} = \frac{dz_2}{\phi_3(x+y)}$$

$$\int \phi_3(x+y) dx = \int dz_2$$

$$x \cdot \phi_3(x+y) + \phi_A(c_A) = z_2$$

$$x \cdot \phi_3(x+y) + \phi_A(x+y) = z_2$$

$$z_2 = x \cdot \phi_3(x+y) + \phi_A(x+y)$$

$$\therefore z = x \phi_1(x-y) + \phi_2(x-y) + x \phi_3(x+y) + \phi_A(x+y)$$

3) Find the soln of the equation $\nabla_1^2 z = e^{-x} \cos y$

i) which reads to zero as $x \rightarrow \infty$

ii) Has the value $\cos y$ when $x=0$.

Soln:-

$$\text{Given } \nabla_1^2 z = e^{-x} \cos y$$

$$\frac{d^2 z}{dx^2} + \frac{d^2 z}{dy^2} = e^{-x} \cos y$$

$$(D^2 + D'^2)z = e^{-x} \cos y \quad \text{--- (1)}$$

To find C.F. :-

$$\text{Now } (D^2 + D'^2)z = 0$$

Assume $z = e^{ax+by}$ is a soln of above eqn.

$$(D^2 + D'^2) e^{ax+by} = 0$$

$$D^2 e^{ax+by} + D'^2 e^{ax+by} = 0$$

$$a^2 e^{ax+by} + b^2 e^{ax+by} = 0$$

$$(a^2 + b^2) e^{ax+by} = 0$$

$$\therefore a^2 + b^2 = 0 \quad \left[\because e^{ax+by} \neq 0 \right]$$

\therefore The Complementary function is

$$\text{C.F.} = e^{ax+by}, \text{ where } a^2 + b^2 = 0.$$

To find P.I. :-

$$\text{P.I.} = \frac{1}{(D + D'^2)} e^{-x} \cos y$$

$$= e^{-x} \frac{1}{(D - 1)^2 + D'^2} \cos y$$

$$= e^{-x} \frac{1}{D^2 + 1 - 2D + D'^2} \cos y \quad \left. \begin{array}{l} \cos(0x + y) \\ a = 0, b = 1 \end{array} \right\}$$

$$= e^{-x} \frac{1}{1 + 1 - 2D - 1^2} \cos y$$

$$= - \frac{e^{-x}}{20} \cos y$$

$$\frac{1}{11} \neq \int$$

$$= - \frac{e^{-x}}{2} \int \cos y \, dx$$

$$P \cdot I = - \frac{e^{-x}}{2} x \cos y$$

$$\therefore z = C \cdot F + P \cdot I$$

$$z = e^{ax+by} - \frac{e^{-x}}{2} x \cos y$$

In general,

$$z = \sum_r A_r e^{a_r x + b_r y} \rightarrow \frac{x e^{-x}}{2} \cos y \quad | \quad a_r^2 + b_r^2 = 0$$

$e^{-x} \rightarrow 0$

↳ (3)

i) since $z \rightarrow 0$ as $x \rightarrow \infty$

\therefore It is possible only when a_r is must

be negative.

$$\text{Let } a_r = -\lambda_r,$$

$$\text{but } a_r^2 + b_r^2 = 0$$

$$\Rightarrow (-\lambda_r)^2 + b_r^2 = 0$$

$$b_r^2 = -\lambda_r^2 \Rightarrow b_r^2 = -\lambda_r^2$$

$$b_r = \sqrt{-\lambda_r^2}$$

$$b_r = \pm i \lambda_r$$

$$(3) \Rightarrow z = \sum_{r=0}^{\infty} A_r e^{\lambda_r x \pm i \lambda_r y} - \frac{x}{2} e^{-x} \cos y \quad (4)$$

Now,

$$B_r \cos(\lambda_r y + \epsilon_r) = B_r [\cos \lambda_r y \cos \epsilon_r - \sin \lambda_r y \sin \epsilon_r]$$

$$= A_r \cos \lambda_r y \pm i A_r \sin \lambda_r y$$

$$\text{where } A_r = B_r \cos \epsilon_r$$

$$\pm i A_r = -B_r \sin \epsilon_r$$

$$= A_r (\cos \lambda_r y \pm i \sin \lambda_r y)$$

$$B_r \cos(\lambda_r y + \epsilon_r) = A_r e^{\pm i \lambda_r y}$$

$$(A) \rightarrow z = \sum_{r=0}^{\infty} B_r e^{-\lambda_r x} \cos(\lambda_r y + \epsilon_r) - \frac{x}{2} e^{-x} \cos y$$

ii) since $z = \cos y$, where $x = 0$.

$$z = \sum_{r=0}^{\infty} B_r e^{-\lambda_r x} \cos(\lambda_r y + \epsilon_r) - \frac{x}{2} e^{-x} \cos y$$

$$\cos y = \sum_{r=0}^{\infty} B_r \cos(\lambda_r y + \epsilon_r)$$

$$\Rightarrow \epsilon_r = 0 \quad \forall r$$

and

$$\lambda_r = \begin{cases} 1 & \text{if } r=0 \\ 0 & \text{if } r \neq 0 \end{cases}$$

$$\Rightarrow B_r = \begin{cases} 1 & \text{if } r=0 \\ 0 & \text{if } r \neq 0 \end{cases}$$

$$z = e^{-x} \cos y - \frac{x e^{-x}}{2} \cos y$$

$$z = e^{-x} \cos y (1 - \frac{x}{2})$$

1) show that $\frac{d^2y}{dx^2} + 2k \frac{dy}{dx} = c^2 \frac{d^2y}{dx^2}$ possesses
 solution of the form

$$\sum_{r=0}^{\infty} C_r e^{-kr} \cos(\alpha_r x + \xi_r) \cos(\omega_r t + \delta_r) \text{ where}$$

$$C_r, \alpha_r, \delta_r, \xi_r \text{ are constant and } \omega_r^2 = \alpha_r^2 c^2 - k^2$$

Soln:- $(D^2 + 2kD - c^2 D'^2)y = 0 \quad \text{--- (1)}$

Let $y = e^{ax+bt}$ be the soln of eqn (1)

$$(D^2 + 2kD - c^2 D'^2) e^{ax+bt} = 0$$

$$b^2 e^{ax+bt} + 2kb e^{ax+bt} - c^2 a^2 e^{ax+bt} = 0$$

$$\Rightarrow b^2 + 2kb - c^2 a^2 = 0 \quad \left\{ \because e^{ax+bt} \neq 0 \right.$$

$$b = \frac{-2k \pm \sqrt{4k^2 - 4(-1)(-a^2 c^2)}}{2}$$

$$= \frac{-2k \pm \sqrt{4k^2 - 4(-1)(-a^2 c^2)}}{2} \quad \text{--- (1)}$$

$$= \frac{-2k \pm \sqrt{4k^2 + 4a^2 c^2}}{2}$$

$$= \frac{-2k \pm 2\sqrt{k^2 + a^2 c^2}}{2}$$

$$b = -k \pm \sqrt{k^2 + a^2 c^2}$$

In general,

$$b_r = -k \pm \sqrt{k^2 + a_r^2 c^2}$$

Let us take $a_r^2 = -\alpha_r^2$

$$a_r = \pm i \alpha_r$$

$$\text{then } b_r = -k \pm \sqrt{k^2 - a_r^2 c^2}$$

$$= -k \pm i \sqrt{\alpha_r^2 c^2 - k^2}$$

$$= -k \pm i \omega_r$$

$$b_r = -k \pm i \omega_r, \text{ where } \omega_r^2 = \alpha_r^2 c^2 - k^2$$

$$\therefore y = e^{a_r x + b_r t}$$

$$= e^{a_r x} \cdot e^{b_r t}$$

$$y = e^{\pm i \alpha_r x - kt} \cdot e^{\pm i \omega_r t}$$

In general,

$$y = \sum_{r=0}^{\infty} C_r e^{-kt} e^{\pm i \alpha_r x} e^{\pm i \omega_r t}$$

$$C e^{\pm i \alpha_r x} e^{\pm i \omega_r t} [A \cos \alpha_r x + B \sin \alpha_r x]$$

$$y = \sum_{r=0}^{\infty} C_r e^{-kt} \cos(\alpha_r x + \xi_r) \cos(\omega_r t + \delta_r)$$

Classification of 2nd order P.D.E

(*)

Def:- Quasi linear P.D.E of 2nd Order.

A second order P.D.E, which is linear w.r. to second order partial derivatives r, s, t is said to be quasi linear P.D.E of 2nd order

Ex:

The equation $Rr + Ss + Tt + f(x, y, z, p, q) = 0$ ①
where $f(x, y, z, p, q)$ need not be linear, is
a quasi linear P.D.E.

Here the co-efficients are R, S, T
may be function of x and y , eqn ① is
said to be

i) Elliptic if $S^2 - 4RT < 0$

ii) parabolic if $S^2 - 4RT = 0$

iii) hyperbolic $S^2 - 4RT > 0$, at a point
(x_0, y_0)

Canonical form of PDE:-

Consider the equation

$$Rr + Ss + Tt + f(x, y, z, p, q) = 0 \quad \text{--- ①}$$

Let us the transformation

$$\xi = \xi(x, y), \quad \eta = \eta(x, y), \quad z = z(\xi, \eta)$$

where ξ & η are continuous differentiable
and the Jacobian

$$J = \frac{\partial(\xi, \eta)}{\partial(x, y)}$$

$$J = \begin{vmatrix} \frac{\partial \xi}{\partial x} & \frac{\partial \xi}{\partial y} \\ \frac{\partial \eta}{\partial x} & \frac{\partial \eta}{\partial y} \end{vmatrix}$$

$$J = \xi_x \eta_y - \xi_y \eta_x \neq 0$$

Now,

$$p = \frac{dz}{dx} = \frac{\partial z}{\partial x} (\xi, \eta)$$

$$= \left(\frac{\partial z}{\partial \xi} \cdot \frac{\partial \xi}{\partial x} + \frac{\partial z}{\partial \eta} \cdot \frac{\partial \eta}{\partial x} \right)$$

$$p = z_{\xi} \xi_x + z_{\eta} \eta_x$$

$$q = \frac{\partial z}{\partial y} = \frac{\partial z}{\partial y} (\xi, \eta)$$

$$q = z_{\xi} \xi_y + z_{\eta} \eta_y$$

$$r = \frac{\partial^2 z}{\partial x^2} = \frac{\partial}{\partial x} (z_{\xi} \xi_x + z_{\eta} \eta_x)$$

$$= z_{\xi} \xi_{xx} + \xi_x \frac{\partial z_{\xi}}{\partial x} + z_{\eta} \eta_{xx} + \eta_x \frac{\partial z_{\eta}}{\partial x}$$

$$= z_{\xi} \xi_{xx} + \xi_x \left[\frac{\partial z_{\xi}}{\partial \xi} \cdot \frac{\partial \xi}{\partial x} + \frac{\partial z_{\xi}}{\partial \eta} \cdot \frac{\partial \eta}{\partial x} \right]$$

$$+ z_{\eta} \eta_{xx} + \eta_x \left[\frac{\partial z_{\eta}}{\partial \xi} \cdot \frac{\partial \xi}{\partial x} + \frac{\partial z_{\eta}}{\partial \eta} \cdot \frac{\partial \eta}{\partial x} \right]$$

$$= z_{\xi} \xi_{xx} + \xi_x [z_{\xi\xi} \xi_x + z_{\xi\eta} \eta_x] + z_{\eta} \eta_{xx} +$$

$$\eta_x [z_{\eta\xi} \xi_x + z_{\eta\eta} \eta_x]$$

$$r = z_{\xi} \xi_{xx} + z_{\xi\xi} \xi_x^2 + z_{\eta\eta} \eta_x^2 + 2 \eta_x \xi_x z_{\xi\eta} + z_{\eta} \eta_{xx}$$

$$t = \frac{\partial^2 z}{\partial y^2} = \frac{\partial}{\partial y} (z_{\xi} \xi_y + z_{\eta} \eta_y)$$

$$= z_{\xi} \xi_{yy} + \xi_y \frac{\partial z_{\xi}}{\partial y} + z_{\eta} \eta_{yy} + \eta_y \frac{\partial z_{\eta}}{\partial y}$$

$$= z_{\xi} \xi_{yy} + \xi_y \left[\frac{dz_{\eta}}{d\xi} \frac{d\xi}{dy} + \frac{dz_{\eta}}{d\eta} \frac{d\eta}{dy} \right] \text{---}$$

$$= z_{\xi} \xi_{yy} + \xi_y [z_{\xi\xi} \xi_y + z_{\xi\eta} \eta_y] + z_{\eta} \eta_{yy} + \eta_y$$

$$[z_{\eta\xi} \xi_y + z_{\eta\eta} \eta_y]$$

$$E = z_{\xi} \xi_{yy} + z_{\xi\xi} \xi_y^2 + z_{\eta\eta} \eta_y^2 + 2\eta_y \xi_y z_{\eta\xi} + z_{\eta} \eta_{yy}$$

$$S = \frac{\partial^2 z}{\partial x \partial y} = \frac{\partial}{\partial y} [z_{\xi} \xi_x + z_{\eta} \eta_x] \text{---}$$

$$= \frac{\partial}{\partial y} \left(\frac{\partial z}{\partial x} \right)$$

$$= \frac{\partial}{\partial y} (z_{\xi} \xi_x) + \frac{\partial}{\partial y} (z_{\eta} \eta_x)$$

$$= \frac{dz_{\xi}}{dy} \xi_x + z_{\xi} \xi_x \xi_y + \frac{dz_{\eta}}{dy} \eta_x + \eta_x \eta_y$$

$$= \left(\frac{dz_{\xi}}{d\xi} \frac{d\xi}{dy} + \frac{dz_{\xi}}{d\eta} \frac{d\eta}{dy} \right) \xi_x + z_{\xi} \xi_x \xi_y + \left(\frac{dz_{\eta}}{d\xi} \frac{d\xi}{dy} + \frac{dz_{\eta}}{d\eta} \frac{d\eta}{dy} \right) \eta_x + z_{\eta} \eta_x \eta_y$$

$$= (z_{\xi\xi} \xi_y + z_{\xi\eta} \eta_y) \xi_x + z_{\xi} \xi_x \xi_y +$$

$$(z_{\eta\xi} \xi_y + z_{\eta\eta} \eta_y) \eta_x + z_{\eta} \eta_x \eta_y$$

$$S = z_{\xi\xi} \xi_x \xi_y + z_{\xi} \xi_x \xi_y + z_{\xi\eta} \xi_x \eta_y + z_{\eta\xi} \eta_x \xi_y + z_{\eta\eta} \eta_x \eta_y + z_{\eta} \eta_x \eta_y$$

sub the value eqn (1)

$$R [z_{\xi\xi} \xi_x^2 + z_{\xi\xi} \xi_{xx} + z_{\eta\eta} \eta_x^2 + z_{\eta\eta} \eta_{xx} + 2z_{\xi\eta} \xi_x \eta_x] +$$

$$T [z_{\xi\xi} \xi_y^2 + z_{\xi\xi} \xi_{yy} + 2z_{\xi\eta} \xi_y \eta_y + z_{\eta\eta} \eta_y^2 + \eta_{yy} z_{\eta\eta}] +$$

$$S [z_{\xi\xi} \xi_x \xi_y + z_{\xi\xi} \xi_x \xi_y + z_{\xi\eta} \xi_x \eta_y + z_{\eta\eta} \eta_x \eta_y + z_{\eta\eta} \eta_x \eta_y +$$

$$z_{\eta\xi} \eta_x \xi_y] + f(x, y, z, p, q) = 0$$

$$\Rightarrow z_{\xi\xi} (R \xi_x^2 + T \xi_y^2 + S \xi_x \xi_y) + z_{\eta\eta} (R \eta_x^2 + T \eta_y^2 + S \eta_x \eta_y)$$

$$+ z_{\eta\xi} (2R \xi_x \eta_x + 2T \xi_y \eta_y + S (\xi_x \eta_y + \eta_x \xi_y)) +$$

$$f(\xi, \eta, z, z_{\xi}, z_{\eta}) = 0$$

$$\therefore A(\xi_x, \xi_y) z_{\xi\xi} + 2B(\xi_x, \xi_y, \eta_x, \eta_y) + A(\eta_x, \eta_y) z_{\eta\eta} =$$

$$- f(\xi, \eta, z, z_{\xi}, z_{\eta}) \quad \text{--- (2)}$$

where $A(\xi_x, \xi_y) = (R \xi_x^2 + T \xi_y^2 + S \xi_x \xi_y)$.

$$A(\eta_x, \eta_y) = (R \eta_x^2 + T \eta_y^2 + S \eta_x \eta_y)$$

$$2B(\xi_x, \xi_y, \eta_x, \eta_y) = (2R \xi_x \eta_x + 2T \xi_y \eta_y + S (\xi_x \eta_y + \eta_x \xi_y))$$

Put $A = A(\xi_x, \xi_y)$

$$B = 2B(\xi_x, \xi_y, \eta_x, \eta_y)$$

$C = A(\eta_x, \eta_y)$

$$\text{Let } B^2 = 4 [B(\xi_n, \xi_y, \eta_n, \eta_y)]$$

from the above eqn we can show that

$$B^2 - 4AC = (S^2 - 4RT)J$$

$$A[B(\xi_x, \xi_y, \eta_x, \eta_y)]^2 - 4A(\xi_x \eta_y - \eta_x \xi_y) = (S^2 - 4RT)J \quad \text{--- (3)}$$

$$\text{where } J = \xi_x \eta_y - \eta_x \xi_y.$$

Find the canonical form of

i) Hyperbolic $(S^2 - 4RT) > 0$

ii) Parabolic $(S^2 - 4RT) = 0$

iii) Elliptic $(S^2 - 4RT) < 0$

Proof :-

Case (i) :- **Hyperbolic**

$$(S^2 - 4RT) > 0$$

Consider the quadratic equation

$$R\lambda^2 + S\lambda + T = 0$$

has real and distinct roots.

Let λ_1 & λ_2 be the real & distinct

roots of $R\lambda^2 + S\lambda + T = 0$.

choose ξ and η

such that

$$\xi_x = \lambda_1 \xi_y \quad \& \quad \eta_x = \lambda_2 \eta_y \quad \text{--- (1)}$$

$$\lambda_1 = \frac{\xi_{xx}}{\xi_{yy}} \quad \& \quad \lambda_2 = \frac{\eta_{xx}}{\eta_{yy}}$$

Consider,

$$\xi_{xx} - \lambda_1 \xi_{yy} = 0$$

by Lagrange's form

$$\frac{dx}{1} = \frac{dy}{-\lambda_1} = \frac{d\xi}{0}$$

(1) & (2) form

$$\frac{dx}{1} = \frac{dy}{-\lambda_1}$$

$$-\lambda_1 dx = dy$$

$$-\lambda_1 = \frac{dy}{dx}$$

$$\frac{dy}{dx} + \lambda_1 = 0$$

$$\Rightarrow P(x, y) = C_1$$

(1) & (3) form

$$\frac{dx}{1} = \frac{d\xi}{0}$$

$$d\xi = 0$$

$$\text{Hing } \xi = \phi_1(c_1)$$

$$= \phi_1(P(x, y))$$

$$\xi = f_1(x, y)$$

||¹⁴ we have

$$\frac{dy}{dx} + \lambda_2 = 0 \text{ and } \eta = f_2(x, y)$$

w.k.T

$$\begin{aligned} A(\xi_x, \xi_y) &= R\xi_x^2 + S\xi_x\xi_y + T\xi_y^2 \\ &= \xi_y^2 \left[R\left(\frac{\xi_x}{\xi_y}\right)^2 + S\frac{\xi_x}{\xi_y} + T \right] \\ &= \xi_y^2 [R\lambda^2 + S\lambda + T] \end{aligned}$$

$$A(\xi_x, \xi_y) = \xi_y^2 (0) = 0 \quad \left[\begin{array}{l} \lambda = 0 \\ R\lambda^2 + S\lambda + T = 0 \end{array} \right] \Rightarrow A=0$$

$$||^{14} A(\eta_x, \eta_y) = 0 \quad \Rightarrow C=0$$

$$\text{Also, } B^2 - 4AC = S^2 - 4RT > 0$$

$$\therefore B^2 - 4AC > 0$$

$$B^2 - 0 > 0$$

$$B > 0$$

c) in canonical form becomes,

$$A(\xi_x, \xi_y) z_{\xi\xi} + 2B(\xi_x, \xi_y, \eta_x, \eta_y) z_{\xi\eta} + C(\eta_x, \eta_y) z_{\eta\eta} + f(\xi, \eta, z, z_\xi, z_\eta) = 0$$

$$2B(\xi_x, \xi_y, \eta_x, \eta_y) z_{\xi\eta} = -f(\xi, \eta, z, z_\xi, z_\eta)$$

$$z_{\xi\eta} = \frac{-f(\xi, \eta, z, z_\xi, z_\eta)}{2B(\xi_x, \xi_y, \eta_x, \eta_y)}$$

$$z_{\xi\eta} = g(\xi, \eta, z, z_\xi, z_\eta)$$

which is canonical form of hyperbolic.

case cii): - $(S^2 - ART) = 0$ parabolic

Consider, $R\lambda^2 + S\lambda + T = 0$

has equal & real roots such that $\lambda_1 = \lambda_2 = \lambda$

choose ξ

such that,

$$\xi_x = \lambda \xi_y$$

$\xi = f_1(x, y)$ (by case ci)

now $A(\xi_x, \xi_y) = R\xi_x^2 + S\xi_x\xi_y + T\xi_y^2$

$$= \xi_y^2 \left[R \left(\frac{\xi_x}{\xi_y} \right)^2 + S \left(\frac{\xi_x}{\xi_y} \right) + T \right]$$

$$= \xi_y^2 (R\lambda^2 + S\lambda + T)$$

$$A(\xi_x, \xi_y) = 0$$

$$\text{Also } B^2 - 4AC = S^2 - 4RT = 0$$

$$B^2 - 4A(0) = 0$$

$$B^2 = 0$$

$$B = 0$$

eqn (3) becomes,

$$0(z_{\xi\xi}) + 0z_{\eta\xi} + Az_{\eta\eta} = -f(\xi, \eta, z, \xi_x, z_{\eta})$$

$$z_{\eta\eta} = g(\xi, \eta, z, z_{\eta})$$

which is canonical form of parabolic

P.D.E.

Case (iii): $S^2 - 4PT < 0$ Elliptic.

Consider $R\lambda^2 + S\lambda + T = 0$ has two imaginary roots λ_1, λ_2

Consider the transformation

$$\xi = \alpha + i\beta, \quad \eta = \alpha - i\beta, \quad \text{where } \alpha, \beta, \text{ are real.}$$

$$\alpha = \frac{1}{2}(\xi + \eta); \quad \beta = -\frac{i}{2}(\xi - \eta)$$

$$\xi = \frac{1}{2} \begin{bmatrix} 1 & i \\ 1 & -i \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \end{bmatrix}$$

$$\alpha_\xi = \frac{1}{2}(1) = \frac{1}{2}; \quad \beta_\xi = -\frac{i}{2}(1) = -\frac{i}{2}$$

$$\beta = -\frac{i}{2}(\xi - \eta)$$

$$\alpha_\eta = \frac{1}{2}(0+1) = \frac{1}{2}; \quad \beta_\eta = -\frac{i}{2}(-1) = \frac{i}{2}$$

$$Z_\xi = Z_\alpha \alpha_\xi + Z_\beta \beta_\xi$$

$$= Z_\alpha \left(\frac{1}{2}\right) + Z_\beta \left(-\frac{i}{2}\right)$$

$$Z_\xi = \frac{1}{2}(Z_\alpha - iZ_\beta)$$

$$Z_{\xi\eta} = \frac{1}{2} \frac{d}{d\eta} (Z_\alpha - iZ_\beta)$$

$$= \frac{1}{2} \left[Z_{\alpha\alpha} \alpha_\eta + Z_{\alpha\beta} \beta_\eta \right] - \frac{i}{2} \left[Z_{\beta\alpha} \alpha_\eta + Z_{\beta\beta} \beta_\eta \right]$$

$$= \frac{1}{2} \left[Z_{\alpha\alpha} \alpha_\eta + Z_{\alpha\beta} \beta_\eta - i Z_{\beta\alpha} \alpha_\eta - i Z_{\beta\beta} \beta_\eta \right]$$

$$= \frac{1}{2} \left[Z_{\alpha\alpha} \left(\frac{1}{2}\right) + Z_{\alpha\beta} \left(\frac{i}{2}\right) - i Z_{\beta\alpha} \left(\frac{1}{2}\right) - i Z_{\beta\beta} \left(\frac{i}{2}\right) \right]$$

$$Z_{\xi\eta} = \frac{1}{2} \left[Z_{\alpha\alpha} \left(\frac{1}{2}\right) + Z_{\beta\beta} \left(\frac{1}{2}\right) \right] = \frac{1}{4} (Z_{\alpha\alpha} + Z_{\beta\beta})$$

proceeding in the similar lines as in case (i):

$$z_{\eta\eta} = g(\xi, \eta, z, z_{\xi}, z_{\eta}) - C_3$$

$$\frac{1}{4} [z_{\alpha\alpha} + z_{\beta\beta}] = g(\xi, \eta, z, z_{\xi}, z_{\eta})$$

$$[z_{\alpha\alpha} + z_{\beta\beta}] = 4 [g(\xi, \eta, z, z_{\xi}, z_{\eta})]$$

$$z_{\alpha\alpha} + z_{\beta\beta} = \phi(\alpha, \beta, z, z_{\alpha}, z_{\beta})$$

which is canonical form of Elliptic P.D.E.

Q. Reduce the P.D.E $y^2 \frac{\partial^2 z}{\partial x^2} - 2xy \frac{\partial^2 z}{\partial x \partial y} + x^2 \frac{\partial^2 z}{\partial y^2}$

Ans.

into canonical form and find the solution.

Soln:-

The given equation is

$$y^2 r - 2xy s + x^2 t = \frac{y^2}{x} \frac{\partial z}{\partial x} + \frac{x^2}{y} \frac{\partial z}{\partial y}$$

Here $R = y^2$, $S = -2xy$, $T = x^2$

$$S^2 - 4RT = (-2xy)^2 - 4y^2 x^2$$

$$S^2 - 4RT = 0$$

\therefore eqn (i) is parabolic.

Now, $R\lambda^2 + S\lambda + T = 0$

$$y^2 \lambda^2 - 2xy \lambda + x^2 = 0$$

$$(y\lambda - x)^2 = 0$$

$$(y\lambda - x)(y\lambda - x) = 0$$

$$\lambda = \frac{x}{y}, \frac{x}{y}$$

choose ξ and η such that $\xi_x = \lambda \xi_y$.

$$\xi_x - \frac{x}{y} \xi_y = 0$$

by Lagrange's form

$$\frac{dx}{1} = \frac{dy}{-x/y} = \frac{d\xi}{0}$$

(1) & (2) nd form

$$\frac{dx}{1} = \frac{dy}{-x/y}$$

$$x dx = -y dy$$

$$\int dx \frac{x^2}{2} = -\frac{y^2}{2} + C_1$$

$$x^2 + y^2 = C_1$$

from (1) & (3)

$$\frac{dx}{1} = \frac{d\xi}{0}$$

$$d\xi = 0$$

Integ

$$\xi = C_2 = x^2 + y^2$$

||| again choose η which is independent of ξ .

$$\text{i.e. } \eta = x^2 - y^2 \text{ and } \xi = x^2 + y^2$$

$$\eta_x = 2x$$

$$\xi_x = 2x$$

$$\eta_y = -2y$$

$$\xi_y = 2y$$

$$\text{Also, } z = z(\xi, \eta)$$

$$\frac{\partial z}{\partial x} = z_\xi \cdot \xi_x + z_\eta \cdot \eta_x$$

$$= z_\xi \cdot 2x + z_\eta \cdot 2x$$

$$\frac{\partial z}{\partial x} = 2x (z_\xi + z_\eta)$$

$$\frac{\partial z}{\partial y} = z_\xi \xi_y + z_\eta \eta_y$$

$$= z_\xi \cdot 2y + z_\eta (-2y)$$

$$\frac{\partial z}{\partial y} = 2y (z_\xi - z_\eta)$$

$2xy [z_{\xi\xi} - z_{\eta\eta}]$

$$\frac{\partial^2 z}{\partial x^2} = \frac{\partial}{\partial x} \left(\frac{\partial z}{\partial x} \right)$$

$$= \frac{\partial}{\partial x} \left(2x (z_\xi + z_\eta) \right)$$

$$= 2 \left[(z_\xi + z_\eta) \cdot 1 + x (z_{\xi\xi} \xi_x + z_{\xi\eta} \eta_x + z_{\eta\xi} \xi_x + z_{\eta\eta} \eta_x) \right]$$

$$= 2 \left[z_\xi + z_\eta + x (z_{\xi\xi} 2x + 2z_{\xi\eta} 2xy + z_{\eta\eta} (2x)) \right]$$

$$\frac{\partial^2 z}{\partial x^2} = 2 \left[z_{\xi\xi} + z_{\eta\eta} + 2x^2 (z_{\xi\xi\xi} + 2z_{\xi\eta\xi} + z_{\eta\xi\xi}) \right]$$

iii

$$\frac{\partial^2 z}{\partial y^2} = \frac{\partial}{\partial y} (2y (z_{\xi} - z_{\eta}))$$

$$\frac{\partial^2 z}{\partial y^2} = 2 (z_{\xi} - z_{\eta} + 2y^2 (z_{\xi\xi\xi} - 2z_{\xi\eta\xi} + z_{\eta\xi\xi}))$$

$$\frac{\partial^2 z}{\partial x \partial y} = \frac{\partial}{\partial x} \left(\frac{\partial z}{\partial y} \right) = \frac{\partial}{\partial x} (2y (z_{\xi} - z_{\eta}))$$

$$= 2y (z_{\xi\xi} \xi_x + z_{\xi\eta\xi} (\xi_x - \eta_x) - z_{\eta\xi\xi} (\xi_x - \eta_x))$$

$$= 2y (z_{\xi\xi} \cdot 2x + z_{\xi\eta\xi} (2x) - z_{\eta\xi\xi} \cdot 2x)$$

$$= 2y (z_{\xi\xi} \cdot 2x - z_{\eta\xi\xi} \cdot 2x)$$

$$\frac{\partial^2 z}{\partial x \partial y} = 4xy (z_{\xi\xi} - z_{\eta\xi\xi})$$

Sub above values in (1)

$$y^2 \left[2(z_{\xi} + z_{\eta} + 2x^2(z_{\xi\xi\xi} + 2z_{\xi\eta\xi} + z_{\eta\xi\xi})) - 2xy(4xy(z_{\xi\xi} - z_{\eta\xi\xi})) \right] + x^2 \left[2(z_{\xi} - z_{\eta} + 2y^2(z_{\xi\xi\xi} - 2z_{\xi\eta\xi} + z_{\eta\xi\xi})) \right]$$

$$= \frac{y^2}{x} (2x(z_{\xi} + z_{\eta})) + \frac{x^2}{y} (2y(z_{\xi} - z_{\eta}))$$

$$\Rightarrow 2y^2 z_{\xi} + 2y^2 z_{\eta} + 4x^2 y^2 z_{\xi\xi\xi} + 4x^2 y^2 z_{\eta\xi\xi} - 8x^2 y^2 z_{\xi\eta\xi}$$

$$+ 8x^2 y^2 z_{\xi\xi\xi} + 8x^2 y^2 z_{\eta\xi\xi} + 2x^2 z_{\xi} - 2x^2 z_{\eta} + 4x^2 y^2 z_{\xi\xi\xi} - 8x^2 y^2 z_{\xi\eta\xi}$$

$$+ 4x^2 y^2 z_{\eta\xi\xi} = 2y^2 z_{\xi} + 2y^2 z_{\eta} + 2x^2 z_{\xi} - 2x^2 z_{\eta}$$

$$16x^2y^2z_{\eta\eta} = 0$$

we can find $\frac{\partial^2 z}{\partial \eta^2} = 0$

$$\frac{\partial}{\partial \eta} \left(\frac{\partial z}{\partial \eta} \right) = 0$$

$$\frac{\partial z}{\partial \eta} = A$$

$$\partial z = A \partial \eta$$

$$z = A\eta + B$$

where A & B are arbitrary function of ξ

$$\therefore z = \eta A(\xi) + B(\xi)$$

$$z = (x^2 - y^2)A(x^2 + y^2) + B(x^2 + y^2)$$

✓
Q2 Reduce the equation

$$(n-1)^2 \frac{\partial^2 z}{\partial x^2} - y^{2n} \frac{\partial^2 z}{\partial y^2} = ny^{2n-1} \frac{\partial z}{\partial y} \text{ to form}$$

canonical form & its general solution.

Soln:-

The given equation is

$$(n-1)^2 x - y^{2n} k = ny^{2n-1} \frac{\partial z}{\partial y} \quad \text{--- (1)}$$

Here,

$$R = (n-1)^2, \quad b = 0, \quad T = -y^{2n}$$

$$S^2 - 4RT = -4(n-1)^2 (-y^{2n})$$

$$= -4 \left[\frac{n^2}{4} + 1 - 2n \right] (-y^{2n})$$

$$= + \frac{2^{2n}}{y^{2n}} - \frac{2^{2n}}{y^{2n}}$$

$$= + 4(n-1)^2 (y^{2n})^{2n} \quad 4^{2n} y^{2n} \left[1n + \frac{1}{n} - 2 \right]$$

$$= 2^2 (n-1)^2 (y^{2n})^{2n}$$

$$= (2(n-1)y^n)^2 > 0$$

\therefore eqn ① is hyperbolic.

Now,

$$R\lambda^2 + S\lambda + T = 0$$

$$(n-1)^2 \lambda^2 - y^{2n} = 0$$

$$(n-1)^2 \lambda^2 = y^{2n}$$

$$\lambda^2 = \frac{y^{2n}}{(n-1)^2}$$

$$\lambda = \pm \frac{y^n}{(n-1)}$$

$$\lambda = \pm \frac{y^n}{(n-1)}$$

$$\lambda_1 = \frac{y^n}{(n-1)}, \quad \lambda_2 = -\frac{y^n}{(n-1)}$$

choose ξ and η

such that,

$$\xi_2 = \lambda_1 \xi_y \quad \text{and} \quad \eta_x = \lambda_2 \eta_y$$

$$\xi_x = \frac{y^n}{(n-1)} \xi_y \quad \text{and} \quad \eta_x = -\frac{y^n}{(n-1)} \eta_y$$

$$\xi_x - \frac{y^n}{(n-1)} \xi_y = 0 \quad \& \quad \eta_x + \frac{y^n}{(n-1)} \eta_y = 0$$

By Lagrange's form

$$\frac{dx}{1} = \frac{dy}{-\frac{y^n}{(n-1)}} = \frac{d\xi}{0}$$

(1) & (2) nd term,

$$\frac{dx}{1} = \frac{dy}{-\frac{y^n}{(n-1)}}$$

$$dx = -(n-1) y^{-n} dy$$

Integ

$$x = \frac{(1-n) y^{-n+1}}{-n+1} + c$$

By Lagrange's form. $x = y^{1-n} + c$. $xc - y^{1-n} = c$

(1) & (3) nd term,

$$\frac{dx}{1} = \frac{d\xi}{0}$$

$$\int 0 dx = \int d\xi$$

Integ

$$\xi = c = x - y^{1-n}$$

$$q = x + y^{1-n}$$

$$\epsilon_x = 1 ; \quad \epsilon_y = -(1-n)y^{-n}$$

$$\eta_x = 1 ; \quad \eta_y = (1-n)y^{-n}$$

constant
 $z_{\epsilon\eta} = 0$
 $(n-1)^2 = 0$

Also,

$$z = z(\epsilon, \eta)$$

$$\frac{\partial z}{\partial x} = z_{\epsilon} \cdot \epsilon_x + z_{\eta} \cdot \eta_x$$

$$= z_{\epsilon} \cdot 1 + z_{\eta} \cdot 1$$

$$\frac{\partial z}{\partial x} = z_{\epsilon} + z_{\eta}$$

$$\frac{\partial z}{\partial y} = z_{\epsilon} \cdot \epsilon_y + z_{\eta} \cdot \eta_y$$

$$= z_{\epsilon} [-(1-n)y^{-n}] + z_{\eta} \cdot (1-n)y^{-n}$$

$$\frac{\partial z}{\partial y} = (1-n)y^{-n} (z_{\eta} - z_{\epsilon})$$

$$\frac{\partial^2 z}{\partial x^2} = \frac{\partial}{\partial x} \left(\frac{\partial z}{\partial x} \right)$$

$$= \frac{\partial}{\partial x} (z_{\epsilon} + z_{\eta})$$

$$= z_{\epsilon\epsilon} \epsilon_x + z_{\epsilon\eta} \eta_x + z_{\eta\epsilon} \epsilon_x + z_{\eta\eta} \eta_x$$

$$= z_{\epsilon\epsilon} (1) + z_{\epsilon\eta} (1) + z_{\eta\epsilon} (1) + z_{\eta\eta} (1)$$

$$= z_{\epsilon\epsilon} + z_{\epsilon\eta} + z_{\eta\epsilon} + z_{\eta\eta}$$

$$\frac{\partial^2 z}{\partial x^2} = z_{\epsilon\epsilon} + 2z_{\epsilon\eta} + z_{\eta\eta}$$

$$\frac{d^2 z}{dy^2} = \frac{d}{dy} \left[(1-n) y^{-n} (z_\eta - z_\xi) \right]$$

$$= (1-n) \left[(z_\eta - z_\xi) (-n) y^{-n+1} + y^{-n} (z_{\eta\xi} \xi y + z_{\eta\eta} \eta y - z_{\xi\eta} \eta y + z_{\xi\xi} \xi y) \right]$$

$$= (1-n) \left[(z_\eta - z_\xi) (-n) y^{-n+1} + y^{-n} \left[z_{\eta\xi} \xi (1-n) y^{-n} + z_{\eta\eta} (1-n) y^{-n} - z_{\xi\eta} (1-n) y^{-n} - z_{\xi\xi} (1-n) y^{-n} \right] \right]$$

$$= (1-n) \left[(-n) y^{-n+1} (z_\eta - z_\xi) + y^{-n} \left[z_{\eta\eta} (1-n) y^{-n} - 2z_{\eta\xi} (1-n) y^{-n} - z_{\xi\xi} (1-n) y^{-n} \right] \right]$$

$$\left[z_{\eta\eta} (1-n) y^{-n} - 2z_{\eta\xi} (1-n) y^{-n} - z_{\xi\xi} (1-n) y^{-n} \right]$$

$$\frac{d^2 z}{dy^2} = (1-n) \left[(-n) y^{-n+1} (z_\eta - z_\xi) \right] +$$

$$y^{-2n} (1-n) \left[z_{\eta\eta} - 2z_{\eta\xi} - z_{\xi\xi} \right].$$

Sub values of eqn (1)

$$(n-1)^2 \left[z_{\xi\xi} + 2z_{\xi\eta} + z_{\eta\eta} \right] - y^{2n} \left\{ (1-n) \left[-ny^{-n+1} \right. \right.$$

$$\left. \left. (z_\eta - z_\xi) + y^{-2n} (1-n) (z_{\eta\eta} - 2z_{\eta\xi} + z_{\xi\xi}) \right] \right\} =$$

$$n y^{2n-1} (1-n) y^{-n} [z_\eta - z_\xi]$$

$$(n-1)^2 [z_{\xi\xi} + 2z_{\xi\eta} + z_{\eta\eta}] + y^{2n} (1-n) y^{-n+1} (z_{\eta} - z_{\xi})$$

$$-(1-n)^2 (z_{\eta\eta} - 2z_{\xi\eta} + z_{\xi\xi}) = n y^{2n-1} (1-n) y^{+n} (z_{\eta} - z_{\xi})$$

$$(n-1)^2 [z_{\xi\xi} + 2z_{\xi\eta} + z_{\eta\eta}] - [(1-n)^2 (z_{\eta\eta} - 2z_{\xi\eta} + z_{\xi\xi})] =$$

$$(n-1)^2 [z_{\xi\xi} + 2z_{\xi\eta} + z_{\eta\eta}] - [(n-1)^2 (1-n)^2 (z_{\eta\eta} - 2z_{\xi\eta} + z_{\xi\xi})] =$$

$$(n-1)^2 [z_{\xi\xi} + 2z_{\xi\eta} + z_{\eta\eta} - z_{\eta\eta} + 2z_{\xi\eta} - z_{\xi\xi}] = 0$$

$$(n-1)^2 A z_{\xi\eta} = 0$$

$$z_{\xi\eta} = 0$$

$$z = f_1(\xi) \eta + f_2(\eta) \quad \frac{\partial^2 z}{\partial \xi \partial \eta} = 0$$

$$\frac{\partial}{\partial \xi} \left(\frac{\partial z}{\partial \eta} \right) = 0$$

$$z = f_1(\xi) \eta + f_2(\eta)$$

$$z = f_1(x-y^{1-n})(x+y^{1-n}) + f_2(x+y^{1-n})$$

Which is a required solution.

ADJOINT OPERATOR:-

Let $LU = \phi$, where L is a differential operator is given by

$$L = a_0(x) \frac{d^n}{dx^n} + a_1(x) \frac{d^{n-1}}{dx^{n-1}} + \dots + a_{n-1}(x)$$

The adjoint operator L^* associated with L is to form the product vLu

$$\text{i.e.} \int_A^B vLu \, dx = [vu]_A^B - \int_A^B u(L^*v) \, dx$$

which is obtained after repeated integration by parts.

Here, L^* is called the adjoint operator of L .

SELF Adjoint Operator:-

If the operator $L = L^*$ then L is called self adjoint operator.

Problem:-

If L is a operator $R \frac{d^2}{dx^2} + S \frac{d^2}{dx dy} + T \frac{d^2}{dy^2}$

$P \frac{d}{dx} + Q \frac{d}{dy} + Z$ and M is a adjoint operator

defined by $MW = \frac{d^2}{dx^2}(RW) + \frac{d^2}{dx dy}(SW) +$

$\frac{d^2}{dy^2}(TW) - \frac{d}{dx}(PW) - \frac{d}{dy}(QW) + ZW$

then show that $\frac{1100}{1000} \times 2.50 = \frac{1200}{1000} \times 3.15 \times 9.00$

$$\iint_S (wLz - zMw) dx dy = \int_c [U \cos(n, x) + V \cos(m, y)] ds$$

where c is closed curve enclosing on a surface S . where $U = RW \frac{dz}{dx} - z \frac{d}{dx} (RW) - z \frac{d}{dy} (SW) + PzW$ where $v = SW \frac{dz}{dx} + TW \frac{dz}{dy} - z \frac{d}{dy} (TW) + QzW$ if $R_x + \frac{1}{2} S_y = P$ and $\frac{1}{2} S_x + T_y = Q$ show that the operator L

is self adjoint

Soln:-

$$wLz = WR \frac{d^2 z}{dx^2} + wS \frac{d^2 z}{dx dy} + wT \frac{d^2 z}{dy^2} +$$

$$wP \frac{dz}{dx} + wQ \frac{dz}{dy} + wZz$$

$$zMW = z \frac{d^2}{dx^2} (RW) + z \frac{d^2}{dx dy} (SW) +$$

$$z \frac{d^2}{dy^2} (TW) - z \frac{d}{dx} (PW) - z \frac{d}{dy} (QW) + zZz$$

$$wLz - zMW = \left(WR \frac{d^2 z}{dx^2} - z \frac{d^2}{dx^2} (RW) \right) + \left(wS \frac{d^2 z}{dx dy} - z \frac{d^2}{dx dy} (SW) \right) + \left(wT \frac{d^2 z}{dy^2} - z \frac{d^2}{dy^2} (TW) \right) + \left(wP \frac{dz}{dx} - z \frac{d}{dx} (PW) \right) + \left(wQ \frac{dz}{dy} - z \frac{d}{dy} (QW) \right)$$

$$= \frac{d}{dx} \left(wR \frac{dz}{dx} - z \frac{d}{dx} R w - z \frac{d(zw)}{dy} + z P w \right)$$

$$+ \frac{d}{dy} \left(wS \frac{dz}{dx} + wT \frac{dz}{dy} - z \frac{d(wT)}{dy} + Q w z \right)$$

$$wLz - zMw = \frac{du}{dx} + \frac{dv}{dy}$$

$$\left[\begin{aligned} & z \frac{d(Pw)}{dy} + Pw \frac{dz}{dy} \\ & = \frac{d}{dy} (z P w) \end{aligned} \right]$$

$$\iint_S (wLz - zMw) dx dy = \iint_S \left(\frac{du}{dx} + \frac{dv}{dy} \right) dx dy$$

$$= \int_C (2u + 2v) ds$$

$$= \int_C (u \cos(n, ox) + v \cos(n, oy)) ds$$

$$\text{Also } \frac{d^2}{dx^2} (Rw) + \frac{d^2}{dx dy} (Sw) + \frac{d^2}{dy^2} (Tw) -$$

$$\frac{d}{dx} (Pw) - \frac{d}{dy} (Qw) + zw$$

$$= R_{xx} w + R_x w_x + R_{xx} w_x + R w_{xx} + S_y w_x +$$

$$\left[\begin{aligned} R_x + \frac{1}{2} S_y &= P \\ 2R_x + S_y &= 2P \\ S_x + 2T_y &= 2Q \end{aligned} \right]$$

$$2R_x + S_y = 2P$$

$$S_x + 2T_y = 2Q$$

$$S_x w_y + S w_{xy} + w S_{xy} + T_{yy} w + T_y w_y$$

$$+ T_y w_y + T w_{yy} - P_x w - P w_x - Q_y w - Q w_y + zw$$

$$= R w_{xx} + w_{xx} (2R_x + S_y) + S w_{xy} + w_y (2T_y + S_x)$$

$$+ T w_{yy} + w (R_{xx} + S_{xy} + T_{yy} - P_x - Q_y + z) - P w_x$$

$$- Q_y w_y$$

$$= R w_{xx} + 2 w_{xx} P + S w_{xy} + 2 w_y Q + T w_{yy} + w z_1 -$$

$$P w_x - Q_y w_y$$

$$\text{where } z_1 = R_{xx} + S_{xy} + T_{yy} - P_x - Q_y + z$$

$$\therefore M W = R \frac{\partial^2 w}{\partial x^2} + S \frac{\partial^2 w}{\partial x \partial y} + T \frac{\partial^2 w}{\partial y^2} + P \frac{\partial w}{\partial x} + Q \frac{\partial w}{\partial y} + w Z,$$

$$M(W) = L(W)$$

$$\therefore M = L$$

$\therefore L$ is self adjoint.

2) Construct an operator adjoint to the Laplace operator and prove that it is self adjoint.

Proof:-

w.k.t

The Laplace operator L is defined by

$$\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

$$L(u) = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}$$

Here,

$$R=1, S=0, T=1, P=0, Q=0 \text{ and } Z=0$$

The adjoint operator is defined by

$$L^* w = \frac{\partial^2}{\partial x^2} (Rw) + \frac{\partial^2}{\partial x \partial y} (Sw) + \frac{\partial^2}{\partial y^2} (Tw) -$$

$$\frac{\partial}{\partial x} (Pw) - \frac{\partial}{\partial y} (Qw) + Zw$$

$$L^* (w) = \frac{\partial^2}{\partial x^2} (w) + \frac{\partial^2}{\partial y^2} (w)$$

$$L^*(w) = \frac{d^2 w}{dx^2} + \frac{d^2 w}{dy^2}$$

$$L^*(w) = L(w)$$

$\therefore L$ is a self adjoint.

3) Determine the adjoint operator of $L(u)$

$L(u) = c^2 u_{xx} - u_{yy}$ then prove that the operator L is self adjoint.

Proof:-

The adjoint operator of

$$c^2 \frac{d^2}{dx^2} \quad \leftarrow \quad \frac{d^2}{dy^2}$$

$$L(u) = c^2 \frac{d^2 u}{dx^2} \quad \leftarrow \quad \frac{d^2 u}{dy^2}$$

Here,

$$R = c^2, S = 0, T = -1, P = 0, Q = 0, Z = 0$$

$$L^*(u) = \frac{d^2}{dx^2} (Rw) + \frac{d^2}{dx dy} (Sw) + \frac{d^2}{dy^2} (Tw) -$$

$$\frac{d}{dx} (Pw) - \frac{d}{dy} (Qw) + Zw$$

$$= \frac{d^2}{dx^2} (c^2 w) + \frac{d^2}{dy^2} (-w)$$

$$= \frac{d^2}{dx^2} c^2 w - \frac{d^2}{dy^2} w.$$

$$L^*(v) = \frac{\partial^2 v}{\partial x^2} c^2 - \frac{\partial^2 v}{\partial y^2}$$

$\therefore L$ is self adjoint.

Riemann Method:-

Consider the linear hyperbolic PDE

$$\frac{\partial^2 z}{\partial x \partial y} + a \frac{\partial z}{\partial x} + b \frac{\partial z}{\partial y} + cz = f(x, y) \quad \text{--- (1)}$$

where a, b, c are function of x, y .

$$\text{Let } L = \frac{\partial^2}{\partial x \partial y} + a \frac{\partial}{\partial x} + b \frac{\partial}{\partial y} + c$$

eqn (1) becomes

$$L(z) = f(x, y)$$

Let M be the adjoint operator of

L which is given by

$$M(w) = \frac{\partial^2 w}{\partial x \partial y} - \frac{\partial}{\partial x} (aw) - \frac{\partial}{\partial y} (bw) + cw$$

Now,

$$wL(z) - zM(w) = \left(w \frac{\partial^2 z}{\partial x \partial y} + aw \frac{\partial z}{\partial x} + bw \frac{\partial z}{\partial y} + wz \right)$$

$$- \left(z \frac{\partial^2 w}{\partial x \partial y} - z \frac{\partial}{\partial x} (aw) - z \frac{\partial}{\partial y} (bw) + zw \right)$$

$$= \left(w \frac{\partial^2 z}{\partial x \partial y} - z \frac{\partial^2 w}{\partial x \partial y} + aw \frac{\partial z}{\partial x} + z \frac{\partial (aw)}{\partial x} + bw \frac{\partial z}{\partial y} + z \frac{\partial (bw)}{\partial y} \right)$$

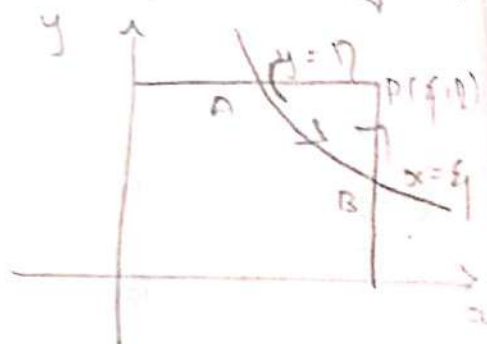
$$= \frac{\partial}{\partial x} \left(wz - z \frac{\partial w}{\partial y} \right) + \frac{\partial}{\partial y} \left(wbz + w \frac{\partial z}{\partial x} \right)$$

$$wz - zmw = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \quad \text{--- (2)} \quad \text{where } u = wz - z \frac{\partial w}{\partial y}$$

$$v = wbz + w \frac{\partial z}{\partial x} \quad \text{--- (3)}$$

Consider the arc of a curve, where the line PA & PB parallel to x & y respectively.

Let 's' denote the area enclosed by the contour ABPA.



Clearly,

$$\text{On AP, } y = \eta \Rightarrow dy = 0$$

$$\text{On PB, } x = \xi \Rightarrow dx = 0$$

Taking

\iint_S on both sides of eqn (2)

$$\iint_S (wz - zmw) dx dy = \iint_S \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) dx dy$$

$$= \int_C (u dy - v dx) = \int_{ABPA} (u dy - v dx)$$

$$\iint_S (wz - zmw) dx dy = \int_A^B (u dy - v dx) + \int_B^P u dy + \int_P^A (-v dx)$$

Now,

$$\text{--- (3)}$$

$$\begin{aligned}
 \int_P^A (v) dx &= \int_P^A (wbz + w \frac{dz}{dx}) dx && \text{L4 (2)} \\
 &= \int_P^A (wbz + w \frac{dz}{dx} + z \frac{dw}{dx} - z \frac{dw}{dx}) dx \\
 &= \int_P^A (wbz - z \frac{dw}{dx}) dx + [wz]_P^A \\
 &= \int_P^A (wbz - z \frac{dw}{dx}) dx + [wz]_A - [wz]_P && \text{L (3A)}
 \end{aligned}$$

Now,

$$(3) \Rightarrow \int_P^A (v) dx = \int_A^B (udy - v dx) + \int_B^P u dy - \iint_S (wbz - zmw) dx dy$$

$$(3A) \Rightarrow [wz]_P^A = \int_P^A (wbz - z \frac{dw}{dx}) dx + (wz)_A - \int_A^B (udy - v dx) - \int_B^P u dy + \iint_S (wbz - zmw) dx dy$$

$$[wz]_P^A = \int_P^A (wbz - z \frac{dw}{dx}) dx + (wz)_A - \int_A^B (udy - v dx - \int_B^P u dy + \iint_S (wbz - zmw) dx dy)$$

$$\begin{aligned}
 [wz]_P^A &= [wz]_A + \int_P^A z (wb - \frac{dw}{dx}) dx - \int_A^B wz (ady - bdx) \\
 &\quad + \int_A^B (z \frac{dw}{dy} dy + w \frac{dz}{dx} dx) - \int_B^P z (wa - \frac{dw}{dy}) dy \\
 &\quad + \iint_S (wbz - zmw) dx dy && \text{--- (4)}
 \end{aligned}$$

Since the function w is arbitrary, we can choose w to satisfy both conditions

i) $M(w) = 0$ on S .

ii) $\frac{\partial w}{\partial y} = aw$ on $x = \xi$

iii) $\frac{\partial w}{\partial x} = bw$ on $y = \eta$

iv) $[w]_P = 1$ (or) $w = 1$ on $P(\xi, \eta)$

$$\textcircled{4} \Rightarrow [z]_P = [wz]_A - \int_A^B wz (ady - bdx) + \int_A^B \left(z \frac{\partial w}{\partial y} dy + w \frac{\partial z}{\partial x} dx \right) + \iint_S wf(x,y) dx dy \quad \textcircled{5}$$

This gives the value of z at any point (ξ, η) when the value of z & $\frac{\partial z}{\partial x}$ are given on the curve AB .

If the value of z and $\frac{\partial z}{\partial y}$ given

then consider

$$\begin{aligned} \int_A^B z \frac{\partial w}{\partial y} dy + w \frac{\partial z}{\partial x} dx &= \int_A^B \left(z \frac{\partial w}{\partial y} + w \frac{\partial z}{\partial y} - w \frac{\partial z}{\partial y} \right) dy \\ &+ \int_A^B \left(w \frac{\partial z}{\partial x} + z \frac{\partial w}{\partial x} - z \frac{\partial w}{\partial x} \right) dx \\ &= \int_A^B \frac{d}{dy} (zw) dy - \int_A^B w \frac{\partial z}{\partial y} dy + \int_A^B \frac{d}{dx} (wz) dx - \int_A^B z \frac{\partial w}{\partial x} dx \\ &= \int_A^B d(zw) - \int_A^B \left(w \frac{\partial z}{\partial y} dy + z \frac{\partial w}{\partial x} dx \right) \end{aligned}$$

$u \frac{\partial v}{\partial x} + v \frac{\partial u}{\partial x} = \frac{\partial}{\partial x} (uv)$
 $\int \frac{\partial}{\partial x} (uv) dx = uv$

$$\int_A^B z \frac{\partial w}{\partial y} dy + w \frac{\partial z}{\partial x} dx = [zw]_B - [zw]_A - \int_A^B (w \frac{\partial z}{\partial y} dy + z \frac{\partial w}{\partial x} dx)$$

Sub (6) in (5)

$$[z]_P = [wz]_A - \int_A^B wz (ady - bdx) + [wz]_B - [wz]_A - \int_A^B (z \frac{\partial w}{\partial x} dx + w \frac{\partial z}{\partial y} dy) + \iint_S wf(x,y) dx dy$$

(5) + (7) =

$$2[z]_P = [wz]_A + [wz]_B - 2 \int_A^B wz (ady - bdx) + \int_A^B [z \frac{\partial w}{\partial y} dy + w \frac{\partial z}{\partial x} dx] - \int_A^B (z \frac{\partial w}{\partial x} dx + w \frac{\partial z}{\partial y} dy) + 2 \iint_S wf(x,y) dx dy$$

$$[z]_P = \frac{[wz]_A + [wz]_B}{2} - \int_A^B wz (ady - bdx) + \frac{1}{2} \int_A^B (z \frac{\partial w}{\partial y} dy + w \frac{\partial z}{\partial x} dx) - \frac{1}{2} \int_A^B (w \frac{\partial z}{\partial y} dy + z \frac{\partial w}{\partial x} dx) + \iint_S wf(x,y) dx dy \quad (8)$$

∴ Eqn (8) is the value of z at any point (ξ, η) when z, $\frac{\partial z}{\partial x}$ and $\frac{\partial z}{\partial y}$ are given on curve r

And eqn (7) is used when z, $\frac{\partial z}{\partial y}$ are given on r

note

Prove that for the equation $\frac{\partial^2 z}{\partial x \partial y} + \frac{1}{4} z = 0$, the

Green's function is given by

$w(x, y, \xi, \eta) = J_0 \sqrt{(x-\xi)(y-\eta)}$ where $J_0(z)$ is Bessel's function on 1st kind with order zero.

Soln:-

Given that $\frac{\partial^2 z}{\partial x \partial y} + \frac{1}{4} z = 0$ — (1)

Here $L = \frac{\partial^2}{\partial x \partial y} + \frac{1}{4}$ then

$a = b = 0$ and $c = \frac{1}{4}$

\therefore adjoint operator of L is self adjoint

$M(w) = \frac{\partial^2 w}{\partial x \partial y} + \frac{1}{4} w$

$$\begin{aligned} wLz - zMw &= \left(w \frac{\partial^2 z}{\partial x \partial y} + w \frac{1}{4} z \right) - \left(z \frac{\partial^2 w}{\partial x \partial y} + z \frac{1}{4} w \right) \\ &= \frac{\partial}{\partial y} \left(w \frac{\partial z}{\partial x} \right) - \frac{\partial}{\partial x} \left(z \frac{\partial w}{\partial y} \right) \end{aligned}$$

$wLz - zMw = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}$, where $u = -z \frac{\partial w}{\partial y}$, $v = w \frac{\partial z}{\partial x}$

$$\iint_S (wLz - zMw) dx dy = \iint_S \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) dx dy$$

$$= \int_C u dy - v dx$$

$$= \int_A^B (u dy - v dx) + \int_B^A (u dy - v dx)$$

$$= \int_A^B (u dy - v dx) - \int_B^P z \frac{\partial w}{\partial y} dy - \int_P^A w \frac{\partial z}{\partial x} dx - \int_P^A z \frac{\partial w}{\partial x} dx$$

$$+ \int_P^A z \frac{\partial w}{\partial x} dx$$

$$= \int_A^B (u dy - v dx) - \int_B^P z \frac{\partial w}{\partial y} dy - \int_P^A \frac{d}{dz} (wz) dx + \int_P^A z \frac{\partial w}{\partial x} dx$$

$$= \int_A^B (u dy - v dx) - \int_B^P z \frac{\partial w}{\partial y} dy - (wz)_A + (wz)_P + \int_P^A z \frac{\partial w}{\partial x} dx$$

$$[wz]_P = - \int_A^B (u dy - v dx) + \int_B^P z \frac{\partial w}{\partial y} dy + [wz]_A - \int_P^A z \frac{\partial w}{\partial x} dx$$

$$+ \iint_S (w_x z - z w_x) dx dy$$

choose w such that,

$$i) w = 1 \text{ at } P(\xi, \eta)$$

$$ii) \frac{\partial w}{\partial y} = 0 \text{ on } x = \xi$$

$$iii) \frac{\partial w}{\partial x} = 0 \text{ on } y = \eta$$

$$iv) M(w) = 0 \text{ on } S$$

$$\text{Let } w = w(f)$$

$$\text{where } f^2 = (x - \xi)(y - \eta)$$

$$2f \frac{df}{dx} = y - \eta$$

$$2f \frac{df}{dy} = x - \xi$$

Assume $w = w(f)$ satisfies condition (i) to (iv)

$$\frac{\partial w}{\partial x} = \frac{dw}{df} \cdot \frac{df}{dx} = \frac{dw}{df} \left(\frac{y - \eta}{2f} \right)$$

$$\frac{\partial^2 w}{\partial x \partial y} = \frac{d}{dy} \left(\frac{\partial w}{\partial f} \left(\frac{y-\eta}{2f} \right) \right)$$

$$= \frac{1}{2} \frac{d}{dy} \left(w_f \cdot f^{-1} (y-\eta) \right)$$

$$= \frac{1}{2} \left[w_{ff} f_y f^{-1} (y-\eta) + w_f (-1) f_y^{-2} (y-\eta) + w_f f^{-1} \right]$$

$$= \frac{1}{2} \left[w_{ff} \left(\frac{x-\xi}{2f} \right) f^{-1} (y-\eta) - w_f f^{-2} \left(\frac{x-\xi}{2f} \right) (y-\eta) + w_f f^{-1} \right]$$

$$= \frac{1}{2} \left[w_{ff} \frac{f^{-1} f^{-2} f^{-1}}{2} - w_f f^{-2} \frac{f^{-1} f^{-2}}{2} + w_f f^{-1} \right]$$

$$= \frac{1}{2} \left[w_{ff} \frac{1}{2} - w_f \frac{f^{-1}}{2} + w_f f^{-1} \right]$$

$$= \frac{1}{4} \left[w_{ff} - w_f f^{-1} + 2w_f f^{-1} \right]$$

$$\frac{\partial^2 w}{\partial x \partial y} = \frac{1}{4} \left[w_{ff} + w_f f^{-1} \right]$$

Using (iv), $w = 0$

$$\therefore M(w) = \frac{1}{4} \left[\frac{\partial^2 w}{\partial f^2} + f^{-1} \frac{\partial w}{\partial f} + w \right] = 0$$

$$4Af^2 \Rightarrow f^2 \frac{\partial^2 w}{\partial f^2} + f \frac{\partial w}{\partial f} + f^2 w = 0 \quad \text{--- (5)}$$

where w is Bessel's function

$$\text{i.e. } w = J_0(f)$$

$$w = J_0 \sqrt{(x-\xi)(y-\eta)} \text{ is a solution}$$

of (2) satisfies (i) to (iv) conditions

$$\text{i) } M(w) = 0 \text{ on } S.$$

$$\text{ii) } \frac{\partial w}{\partial y} = aw \text{ on } x = \xi$$

$$\text{iii) } \frac{d\omega}{dx} = b\omega \text{ on } y = \eta$$

$$\text{iv) } [\omega]_p = 1 \text{ (or) } \omega = 1 \text{ on } p(\xi, \eta)$$

\therefore Putting the value of ω in (2) we get the value of Z . of given p.d.e.

24.07.19

UNIT-II

Elliptic Differential Equation

Derivation of Laplace equation:-

Let two particles of masses m_1 and m_2 be situated at any points P and Q with a distance 'r'.

By Newton's gravitational law

$$F = G \frac{m_1 m_2}{r^2}$$

where G is gravitational constant.

Let us define a position vector

$$\vec{r} = P \rightarrow Q$$

Assume that, unit mass at Q & G.

i.e) $m_2 = 1$ & $G = 1$

$$\therefore F = \frac{m_1}{r^2}$$

The force at Q due to the mass at P is given by

$$\begin{aligned} \vec{F} &= F \cdot \vec{r} \\ &= F \left(\frac{\vec{r}}{r} \right) \\ &= \frac{F}{r} \left(\vec{r} \right) \\ &= m_1 \nabla \left(\frac{1}{r} \right) \end{aligned}$$

$$\vec{F} = \nabla \left(\frac{m_1}{r} \right) \quad \text{--- (1)}$$



$$\nabla = \sum \frac{\partial}{\partial x_i} \left(\frac{1}{r} \right)$$

$$\nabla \left(\frac{1}{r} \right) = - \frac{\vec{r}}{r^3}$$

$$\begin{aligned} r^2 &= x^2 + y^2 + z^2 \\ r &= \sqrt{x^2 + y^2 + z^2} \\ |\vec{r}| &= \sqrt{x^2 + y^2 + z^2} \end{aligned}$$

where \vec{F} is called the intensity of gravitational force.

Suppose a particle of the unit mass moves under a attraction of the particle of m_1 at P from ∞ to R .

Then the workdone is

$$W = \int_{\infty}^R \vec{F} \cdot d\vec{r}$$

$$= \int_{\infty}^R \nabla \left(\frac{m_1}{r} \right) dr$$

$$= \int_{\infty}^R \left(\sum \frac{d}{dx} \left(\frac{m_1}{r} \right) \vec{i} \right) \cdot (\sum dx \vec{i})$$

$$= \int_{\infty}^R \sum \frac{d}{dx} \left(\frac{m_1}{r} \right) dx$$

$$= \int_{\infty}^R d \left(\frac{m_1}{r} \right)$$

$$= \left(\frac{m_1}{r} \right)_{\infty}^R \Rightarrow \frac{m_1}{r} - 0$$

$$W = \frac{m_1}{r}$$

This defines a potential energy ϕ due the particle of P is denoted by

$$V = - \frac{m_1}{r} \quad \text{--- (2)} \quad \phi = -W$$

from (1)

$$\Rightarrow \vec{F} = \nabla \left(\frac{m_1}{r} \right)$$

$$\vec{F} = \nabla(-V) \quad \text{--- (3)}$$

Consider the system of masses m_1, m_2, \dots, m_n which are h-distance r_1, \dots, r_n respectively.

Then the force of attraction for unit mass due to the system of masses

$$m_1, m_2, \dots, m_n \quad \text{is} \quad \vec{F} = \vec{F}_1 + \vec{F}_2 + \dots + \vec{F}_n$$

$$\vec{F} = \nabla \left(\frac{m_1}{r_1} \right) + \nabla \left(\frac{m_2}{r_2} \right) + \dots + \nabla \left(\frac{m_n}{r_n} \right)$$

$$= \sum_{i=1}^n \nabla \left(\frac{m_i}{r_i} \right)$$

$$\vec{F} = \nabla \left(\sum_{i=1}^n \frac{m_i}{r_i} \right)$$

Then work done is

$$W = \int_A^B \vec{F} \cdot d\vec{r}$$

$$= \int_{\infty}^r \vec{F} \cdot d\vec{r}$$

$$= \int_{\infty}^r \nabla \left(\sum_{i=1}^n \frac{m_i}{r_i} \right) d\vec{r}$$

$$= \int_{\infty}^r \left[\sum \frac{\partial}{\partial x} \left(\sum_{i=1}^n \frac{m_i}{r_i} \right) \vec{i} \right] \cdot \sum dx \vec{i}$$

$$\left(\frac{1}{r} \right) \frac{\partial}{\partial x} \frac{1}{r} = -\frac{1}{r^2} \frac{x}{r} = -\frac{x}{r^3}$$

$$= \int_0^r \sum_{i=1}^n \frac{d}{dx} \left(\frac{m_i}{r_i} \right) dx$$

$$= \int_0^r d \left(\sum_{i=1}^n \frac{m_i}{r_i} \right)$$

$$= \left[\sum_{i=1}^n \frac{m_i}{r_i} \right]_0^r$$

$$= \sum_{i=1}^n \frac{m_i}{r_i}$$

$$W = -V$$

$$\therefore V = - \sum_{i=1}^n \frac{m_i}{r_i}$$

$$\therefore \nabla^2 V = \nabla^2 \left(- \sum_{i=1}^n \frac{m_i}{r_i} \right)$$

$$= - \sum_{i=1}^n \nabla^2 \left(\frac{m_i}{r_i} \right)$$

$$\nabla^2 V = 0 \quad (\text{Laplace equation}) \quad \left[\because \nabla^2 \left(\frac{m_i}{r_i} \right) = 0 \right]$$

where $\nabla^2 = \nabla \cdot \nabla$

$$= \left(\vec{i} \frac{d}{dx} + \vec{j} \frac{d}{dy} + \vec{k} \frac{d}{dz} \right) \cdot \left(\vec{i} \frac{d}{dx} + \vec{j} \frac{d}{dy} + \vec{k} \frac{d}{dz} \right)$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

This is called Laplace equation.

In case the matter of density ρ is continuously distributed in a volume V .

$$V = \iiint_V \rho(x, y, z) dV$$

are r is calculated by

$$r = \sqrt{(x-\xi)^2 + (y-\eta)^2 + (z-\zeta)^2}$$

Def:- Gauss Law:-

$$\iint_S \vec{g} \cdot d\vec{s} = -4\pi G M$$

where \vec{g} = gravitational field ^{on} of system.

M = Total mass inside the system.

G = Gravitational Constant.

26.7.19
20
Derivation of Poisson's equation:-

Let S be a closed surface consisting of particles of masses

m_1, m_2, \dots, m_n .



Let Q be any point on S and

Let $\sum_{i=1}^n m_i = M$ be the total mass inside

S .

Let $\vec{g}_1, \dots, \vec{g}_n$ be the gravitational

field at Q due to m_1, m_2, \dots, m_n respectively

in S .

Also,

$$\vec{g} = \sum_{i=1}^n \vec{g}_i$$

By Gauss law,

$$\begin{aligned} \iint_S \vec{g} \cdot d\vec{s} &= -4\pi G M \\ &= -4\pi G (\rho V) \\ &= -4\pi G \iiint_V \rho dV \quad \text{--- (1)} \end{aligned}$$

$\rho = \frac{\text{mass}}{\text{Volume}}$
 $m = \rho V$

where $M = \iiint_V \rho dV$

ρ = The mass density function.

Since the gravitational field is a conservative field.

i.e) $\vec{g} = \nabla V$

where V is the scalar potential energy.

But, the Gauss divergence theorem

$$\iint_S \vec{g} \cdot d\vec{s} = \iiint_V \nabla \cdot \vec{g} \cdot dV \quad \text{--- (2)}$$

From (1) - (2)

$$0 = -4\pi G \iiint_V \rho dV - \iiint_V \nabla \cdot \vec{g} \cdot dV$$

$$= \iiint_V (-4\pi G \rho - \nabla \cdot \vec{g}) dV$$

$$\therefore -4\pi G \rho - \nabla \cdot \vec{g} = 0$$

$$\nabla \cdot \vec{g} = -4\pi G \rho$$

$$\nabla \cdot (\nabla V) = -4\pi G \rho$$

$$\nabla^2 V = -4\pi G \rho$$


which is the Poisson Equation.

Boundary value problem for Laplace equation..

Def:-

Boundary value problem for i^{th} kind: (or)

1st Dirichlet Problem ^{continuous} Dirichlet Problem

If $f \in C^{(0)}$ and is prescribed on the boundary, C of a finite region R . 

The problem is determined the function $\phi(x, y, z)$ such that,

$$\nabla^2 \phi = 0 \text{ in } R \text{ and satisfying } \phi = f \text{ on } C$$

is called the i^{th} kind boundary value problem (or) Dirichlet problem.

2nd kind (or) Neumann problem:-

This 2nd kind of boundary value problem is determined by $\phi(x, y, z)$ such that

$$\nabla^2 \phi = 0 \text{ with in } R \text{ while } \frac{\partial \phi}{\partial x} \text{ is satisfied}$$

at every point on 'C'.

where $\frac{\partial \phi}{\partial x}$ is the normal derivative

of ϕ , then above problem is 2nd kind (or)

Neumann problem.

3rd kind (or) Mixed boundary value problem
(or) churchills problem

This problem is determine the function $\phi(x, y, z)$ such that $\nabla^2 \phi = 0$ with in R, while the boundary condition $\frac{d\phi}{dx} + h\phi = f$ on 'c' when $h \geq 0$, is

satisfied at every point on the boundary 'c'.

This problem is called 3rd kind (or) mixed boundary value problem (or)

churchills problem. \rightarrow 3. Separation of variables method

Find the solution of the 2-D laplace eqn by using separation of variable method.

$$\nabla^2 U = \frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} = 0$$

of variable method.

Soln:- Consider the 2-D laplace equation

$$\nabla^2 U = \frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} = 0 \quad \text{--- (1)}$$

Assume that the solution of the eqn (1) is of the form

$$U(x, y) = X(x) \cdot Y(y) \quad \text{--- (2)}$$

$$\text{(1)} \Rightarrow \frac{\partial^2}{\partial x^2} (X(x) Y(y)) + \frac{\partial^2}{\partial y^2} (X(x) Y(y)) = 0$$

$$Y(y) \frac{\partial^2 X}{\partial x^2} + X(x) \frac{\partial^2 Y}{\partial y^2} = 0$$

$$Y \frac{\partial^2 X}{\partial x^2} = -X \frac{\partial^2 Y}{\partial y^2}$$

$$\frac{1}{x} \frac{d^2 x}{dx^2} = -\frac{1}{y} \frac{d^2 y}{dy^2} = k$$

where k is a separation parameter.

Now,

$$\frac{d^2 x}{dx^2} = xk, \quad \frac{d^2 y}{dy^2} = -ky$$

$$\frac{d^2 x}{dx^2} - kx = 0, \quad \frac{d^2 y}{dy^2} + ky = 0 \quad \text{--- (3)}$$

Case (i):-

if $k > 0$

Let $k = p^2$, p is real

$$(3) \quad \frac{d^2 x}{dx^2} - p^2 x = 0$$

$$(D^2 - p^2)x = 0$$

$$m^2 - p^2 = 0$$

$$m^2 = p^2$$

$$m = \pm p$$

$$\therefore \text{C.F. is } x(x) = C_1 e^{px} + C_2 e^{-px}$$

Next $k = p^2$ in (3)

$$\frac{d^2 y}{dy^2} + p^2 y = 0$$

$$D^2 + p^2 = 0$$

$$m^2 + p^2 = 0$$

$$m^2 = -p^2$$

$$m = \pm ip$$

$$\therefore \text{C.F. is } y(y) = C_3 \cos py + C_4 \sin py$$

$$\therefore U(x, y) = (C_1 e^{px} + C_2 e^{-px}) (C_3 \cos py + C_4 \sin py)$$

(Casecii): $-k < 0$

$$\text{Let } k = -p^2$$

$$(3) \Rightarrow \frac{d^2x}{dx^2} - (-p^2)x = 0$$

$$\frac{d^2x}{dx^2} + p^2x = 0$$

$$D^2 + P^2 x = 0$$

$$m^2 + p^2 = 0$$

$$m^2 = -p^2$$

$$m = \pm ip$$

$$\therefore \text{C.F. is } x(x) = c_1 \cos px + c_2 \sin px.$$

Next, $k = -p^2$ in (3)

$$\frac{d^2y}{dy^2} + (-p^2)y = 0$$

$$(D^2 - p^2)y = 0$$

$$m^2 - p^2 = 0$$

$$m^2 = p^2$$

$$m = \pm p$$

$$\therefore \text{C.F. is } y(y) = c_3 e^{py} + c_4 e^{-py}$$

$$\therefore U(x, y) = (c_1 \cos px + c_2 \sin px) \cdot (c_3 e^{py} + c_4 e^{-py})$$

(Caseciii) $k = 0$

$$(3) \Rightarrow \frac{d^2x}{dx^2} = 0 ; \frac{d^2y}{dy^2} = 0$$

$$\text{Integ} \Rightarrow \int \frac{d}{dx} \left(\frac{dx}{dx} \right) dx = 0 \quad , \quad \int \frac{d}{dy} \left(\frac{dy}{dy} \right) dy = 0$$

$$\therefore \frac{dx}{dx} = c_1 \quad , \quad \frac{dy}{dy} = c_2$$

ling $\therefore x = x_1 + c_2, y = y_1 c_3 + c_4$

$\therefore u(x, y) = (x_1 + c_2) \cdot (y_1 c_3 + c_4)$

Remark:-

Polar Co-Ordinates:-

The Polar Co-Ordinates (r, θ) is defined by $x = r \cos \theta, y = r \sin \theta$ and also $r^2 = x^2 + y^2, \theta = \tan^{-1}(y/x)$

✓ Cylindrical Co-Ordinates

Cylindrical Co-Ordinates (r, θ, z) is

defined by $x = r \cos \theta, y = r \sin \theta, z = z$ and

also we have $r^2 = x^2 + y^2, \theta = \tan^{-1}(y/x),$

$z = z.$

Spherical Co-Ordinates:-

(r, θ, ϕ) is defined by $x = r \cos \theta \cos \phi,$

$y = r \sin \theta \sin \phi, z = r \cos \theta$ and also $r^2 = x^2 + y^2 + z^2,$

$$\theta = \cos^{-1} \frac{z}{\sqrt{x^2 + y^2 + z^2}}, \phi = \tan^{-1}(y/x)$$

Remark:-

i) Laplace Equation in polar Co-Ordinates

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = 0$$



In Cylindrical Co-Ordinates

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} + \frac{\partial^2 u}{\partial z^2} = 0$$

(iii) Spherical Co-Ordinates:-

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \cdot \frac{\partial u}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 u}{\partial \phi^2} = 0$$

Problem:-

show that if the 2-D Laplace equation $\nabla^2 u = 0$ is transformed by plane polar co-ordinates (r, θ) is defined by the relation $x = r \cos \theta$, $y = r \sin \theta$ it takes

$$\text{the form } \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = 0$$

The Laplace equation in Cartesian co-ordinates is

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0 \quad \text{--- (1)}$$

The Relation between Cartesian and Polar - coordinates are

$$x = r \cos \theta, \quad y = r \sin \theta \quad \text{--- (2)}$$

$$\text{The relation is } r^2 = x^2 + y^2, \quad \theta = \tan^{-1} \left(\frac{y}{x} \right) \quad \text{--- (3)}$$

Diff eqn (1) with respect to x & y .

$$2r \frac{dr}{dx} = 2x, \quad 2r \frac{dr}{dy} = 2y$$

$$r_{xc} = \frac{\partial c}{\partial x} = r_y = \frac{y}{r}$$

$$r_{xc} = \frac{r \cos \theta}{r} = \cos \theta \quad r_y = \frac{r \sin \theta}{r} = \sin \theta \quad r_z = 0$$

(3) diff w.r. to x & y

$$\theta_{xc} = \frac{1}{1+y^2/x^2} y \left(-\frac{1}{x^2} \right)$$

$$= -\frac{y}{x^2+y^2}$$

$$= -\frac{r \sin \theta}{r^2 (\sin^2 \theta + \cos^2 \theta)}$$

$$\theta_{xc} = -\frac{\sin \theta}{r}$$

since

$$\theta_{yc} = \frac{1}{1+y^2/x^2} \left(\frac{1}{x} \right)$$

$$= \frac{x}{x^2+y^2}$$

$$= \frac{r \cos \theta}{r^2}$$

$$\theta_{yc} = \frac{\cos \theta}{r}$$

$$\theta_z = 0$$

Since $U = U(r, \theta)$

$$U_x = U_r r_{xc} + U_\theta \theta_{xc} + U_z r_{zc}$$

$$U_{xc} = U_r \cos \theta + U_\theta \left(-\frac{\sin \theta}{r} \right) \quad \text{--- (4)}$$

$$U_y = U_r \sin \theta + U_\theta \left(\frac{\cos \theta}{r} \right) \quad \text{--- (5)}$$

(A) diff w.r.t to x

$$U_{xx} = U_{xrr} r_x + U_{x\theta} \theta_x \quad \text{--- (6)}$$

To find U_{xr}

$$(A) \Rightarrow U_x = U_r \cos\theta - U_\theta \frac{\sin\theta}{r}$$

$$U_{xr} = \cos\theta U_{rr} - \sin\theta \left(\frac{1}{r} U_{\theta r} + U_\theta \left(-\frac{1}{r^2}\right) \right) \quad \text{--- (7)}$$

To find $U_{x\theta}$

$$(A) \Rightarrow U_{x\theta} = U_r \cos\theta + U_\theta (-\sin\theta) - \frac{1}{r} (U_{\theta\theta} \sin\theta + U_\theta \cos\theta)$$

$$U_{x\theta} = U_r \cos\theta - U_\theta \sin\theta - \frac{U_{\theta\theta} \sin\theta}{r} - \frac{U_\theta \cos\theta}{r} \quad \text{--- (8)}$$

Sub (7) & (8) in (6)

$$U_{xx} = \left[\cos\theta U_{rr} - \sin\theta \left(\frac{1}{r} U_{\theta r} + U_\theta \left(-\frac{1}{r^2}\right) \right) \right] \cos\theta + \left[U_r \cos\theta - U_\theta \sin\theta - \frac{U_{\theta\theta} \sin\theta}{r} - \frac{U_\theta \cos\theta}{r} \right] \left(-\frac{\sin\theta}{r} \right)$$

$$U_{xx} = U_{rr} \cos^2\theta - U_{\theta r} \frac{\sin\theta \cos\theta}{r} + U_\theta \frac{\sin\theta \cos\theta}{r^2} - U_{\theta r} \frac{\sin\theta \cos\theta}{r} + U_r \frac{\sin^2\theta}{r} + U_{\theta\theta} \frac{\sin^2\theta}{r^2} + U_\theta \frac{\sin\theta \cos\theta}{r^2} \quad \text{--- (A)}$$

||| y

$$U_{yy} = U_{rr} \sin^2\theta + U_{\theta r} \frac{\sin\theta \cos\theta}{r} - U_\theta \frac{\sin\theta \cos\theta}{r^2}$$

$$+ U_{\theta r} \frac{\sin\theta \cos\theta}{r} + U_r \frac{\cos^2\theta}{r} + U_{\theta\theta} \frac{\cos^2\theta}{r^2} - U_\theta \frac{\sin\theta \cos\theta}{r^2} \quad \text{--- (B)}$$

$$(A) + (B) = z$$

$$U_{xx} + U_{yy} = U_{rr} + \frac{U_r}{r} + \frac{U_{\theta\theta}}{r^2} = 0$$

$$\therefore \frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} + \frac{1}{r^2} \frac{\partial^2 U}{\partial \theta^2} = 0$$

2) Show that in cylindrical co-ordinates (r, θ, z) defined by the relation $x = r \cos \theta$, $y = r \sin \theta$, $z = z$ the Laplace equation $\nabla^2 U = 0$ takes the form

$$\frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} + \frac{1}{r^2} \frac{\partial^2 U}{\partial \theta^2} + \frac{\partial^2 U}{\partial z^2} = 0$$

Soln:- 1st Problem Proof \rightarrow Just of one change

Result:-

$$i) \underline{x^2 R'' + x R' + (m^2 x^2 - n^2) R = 0}$$
 is a Bessel's equation of order 'n', where n is an integer. Then the solution is

$R(x) = A_m J_n(mx) + B_m Y_n(mx)$

where J_n and Y_n are Bessel function of 1st and 2nd order respectively. Also know that $Y_n \rightarrow \infty$ as $x \rightarrow 0$.

ii) $\underline{x^2 R'' + x R' + (m^2 x^2) R = 0}$ is a Bessel equation of order 'zero'.

The solution is $\underline{R(x) = A_m J_0(mx) + B_m Y_0(mx)}$

where $Y_0 \rightarrow \infty$ as $x \rightarrow 0$

$$\text{iii) } x^2 R'' + xR' + (m^2 x^2 - n^2)R = 0 \text{ where } n$$

is not an integer.

The solution is

$$R(x) = A_{mn} J_n(mx) + B_{mn} J_{(-n)}$$

Bessel function of 1st kind and also J_{-n}

as $x \rightarrow 0$

iv) $x^2 R'' + xR' - n^2 R = 0$ is a homogeneous

linear equation. Then the solution is

$$R(x) = A_n x^n + B_n x^{-n}$$

1) Find the soln of Laplace equation in cylindrical co-ordinates (or) separation of variable method.

soln:-

The Laplace equation in cylindrical are (r, θ, z)

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} + \frac{\partial^2 u}{\partial z^2} = 0 \quad \text{--- (1)}$$

Let $u(r, \theta, z) = R(r) \theta(\theta) z(z)$ is the solution of equation (1)

From (1)

$$\Rightarrow \frac{\partial^2}{\partial r^2} [R(r) \theta(\theta) z(z)] + \frac{1}{r} \frac{\partial}{\partial r} [R(r) \theta(\theta) z(z)] +$$

$$\frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} [R(r) \theta(\theta) z(z)] + \frac{\partial^2}{\partial z^2} [R(r) \theta(\theta) z(z)] = 0$$

$$\Rightarrow \theta z R'' + \frac{1}{r} \theta z R' + \frac{\theta''}{r^2} z R + R \theta z'' = 0$$

by $R\theta z$

$$\frac{R''}{R} + \frac{1}{r} \frac{R'}{R} + \frac{1}{r^2} \frac{\theta''}{\theta} + \frac{z''}{z} = 0 \quad \text{--- (2)}$$

$$\Rightarrow \frac{z''}{z} = -\frac{R''}{R} - \frac{1}{r} \frac{R'}{R} - \frac{1}{r^2} \frac{\theta''}{\theta} = m^2 \text{ (say)}$$

$$\Rightarrow \frac{z''}{z} = m^2 \text{ (say)} \quad \text{--- (3)}$$

Sub eqn (3) in eqn (2) we have

$$\frac{R''}{R} + \frac{1}{r} \frac{R'}{R} + \frac{1}{r^2} \frac{\theta''}{\theta} + m^2 = 0 \quad \text{--- (4)}$$

$$\text{(3) } \Rightarrow z'' - m^2 z = 0$$

i.e. $(D^2 - m^2)z = 0$

Auxiliary equation

$$p^2 - m^2 = 0$$

$$p^2 = m^2$$

$$p = \pm m$$

Complementary function is

$$z(z) = C_1 e^{mz} + C_2 e^{-mz}$$

from (4) \Rightarrow

$$\frac{R''}{R} + \frac{1}{r} \frac{R'}{R} + \frac{1}{r^2} \frac{\theta''}{\theta} + m^2 = 0$$

by r^2

$$r^2 \frac{R''}{R} + r \frac{R'}{R} + \frac{\theta''}{\theta} + m^2 r^2 = 0$$

$$\Rightarrow r^2 \frac{R''}{R} + r \frac{R'}{R} + m^2 r^2 = -\frac{\theta''}{\theta} = n^2 \text{ (say)}$$

$$\text{i.e. } r^2 R'' + r R' + m^2 r^2 R = n^2 R \quad \text{--- (5)}$$

$$\text{and } \theta'' + n^2 \theta = 0 \quad \text{--- (6)}$$

$$\text{(5)} \Rightarrow r^2 R'' + r R' + (m^2 r^2 - n^2) R = 0$$

which is a Bessel equation of Order 'n'.

If 'n' is the integer then the solution is

$$R(r) = C_3 J_n(mr) + C_4 Y_n(mr)$$

$$\text{eqn (6)} \Rightarrow \theta'' + n^2 \theta = 0$$

$$\text{i.e. } (D^2 + n^2) \theta = 0$$

Auxiliary equation is $\theta^2 + n^2 = 0$

$$\Rightarrow \theta = \pm in$$

Complementary function is

$$\theta(\theta) = C_5 \cos n\theta + C_6 \sin n\theta$$

$$\therefore u(r, \theta, z) = R(r) \theta(\theta) z(z)$$

$$u = [C_1 e^{mz} + C_2 e^{-mz}] [C_3 J_n(mr) + C_4 Y_n(mr)] [C_5 \cos n\theta + C_6 \sin n\theta]$$

Since $Y_n(mr) \rightarrow \infty$ as $r \rightarrow 0$

\therefore This solution u is unbounded at $r=0$
 In this case we can choose $C_4 = 0$.

$$u = C_3 J_n(mr) (C_1 e^{mz} + C_2 e^{-mz}) (C_5 \cos n\theta + C_6 \sin n\theta)$$

$$u = J_n(mr) (A \cos n\theta + B \sin n\theta) (C_1 e^{mz} + C_2 e^{-mz})$$

General soln of Bessel eqn of order 'n'
 $y(x) = A J_n(x) + B Y_n(x)$

Also $z \rightarrow 0$ as $-z \rightarrow 0$
 function this case, we can choose $c_1 = 0$.

$\therefore u = J_n(mr) e^{-mz} (A' \cos n\theta + B' \sin n\theta)$
 which is the required solution.

RESULT:-

$$i) \theta'' + \cot \theta \theta' + \left(n(n+1) - \frac{m^2}{\sin^2 \theta} \right) \theta = 0$$

$$(1-\mu^2)^2 \frac{d^2 \theta}{d\mu^2} - 2\mu \frac{d\theta}{d\mu} + \left(n(n+1) - \frac{m^2}{1-\mu^2} \right) \theta = 0$$

where $\mu = \cos \theta$.

This is Legendre's assumption equation then the solution becomes.

$$\theta(\theta) = A_{nm} P_n^m(\cos \theta) + B_{nm} Q_n^m(\cos \theta)$$

where P_n^m and Q_n^m are associated Legendre solution of 1st and 2nd kind respectively.

Note that $Q_n^m \rightarrow \infty$ as $\theta \rightarrow 0$.

$$ii) \theta'' + \cot \theta \theta' + n(n+1)\theta = 0$$

$$(or) (1-\mu^2) \frac{d^2 \theta}{d\mu^2} - 2\mu \frac{d\theta}{d\mu} + n(n+1)\theta = 0$$

where, $\mu = \cos \theta$.

which is a Legendre equation whose solution is

$$\theta(\theta) = A_n P_n(\cos \theta) + B_n Q_n(\cos \theta)$$

(or)

$$\theta(\mu) = A_n P_n(\mu) + B_n Q_n(\mu)$$

where P_n and Q_n are Legendre functions of 1st and 2nd kind.

Now $Q_n \rightarrow \infty$ as $\theta \rightarrow 0$.

$$\text{iii) } r^2 R'' + 2rR' + [\lambda^2 r^2 - n(n+1)]R = 0$$

where solution is

$$R(r) = \lambda^{-1/2} \left[(A_n)(n+1/2) J_{n+1/2}(\lambda r) + B_n J_{-n+1/2}(\lambda r) \right]$$

where,

$J_{\pm n+1/2}$ are spherical Bessel functions.

Q. Solve the Laplace equation in spherical coordinates by using the method of separation of variables.

The Laplace equation in spherical coordinates (r, θ, ϕ) is

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial u}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 u}{\partial \phi^2} = 0$$

Multiplying by r^2 on both sides,

$$\frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial u}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 u}{\partial \phi^2} = 0 \quad \text{--- (1)}$$

Let $u(r, \theta, \phi) = R(r)\Theta(\theta)\Phi(\phi)$ be a solution of (1)

From (1) & (2) \Rightarrow

$$\frac{\partial}{\partial r} \left[r^2 \frac{\partial}{\partial r} (R(r)\Theta(\theta)\Phi(\phi)) \right] + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left[\sin \theta \frac{\partial}{\partial \theta} (R(r)\Theta(\theta)\Phi(\phi)) \right] + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} (R(r)\Theta(\theta)\Phi(\phi)) = 0$$

$$\Rightarrow \frac{d}{dr} [r^2 R' \phi \theta] + \frac{1}{\sin \theta} \frac{d}{d\theta} [\sin \theta R \phi \theta'] +$$

$$\frac{1}{\sin^2 \theta} R \theta \phi'' = 0$$

$$\Rightarrow r^2 \theta \phi R'' + 2r \theta \phi R' + \frac{R \phi}{\sin \theta} (\sin \theta \theta'') + \frac{1}{\sin \theta} R \phi \theta' \cos \theta$$

$$+ \frac{1}{\sin^2 \theta} R \theta \phi'' = 0$$

$$\Rightarrow r^2 \theta \phi R'' + \theta \phi R' (2r) + R \phi \theta'' + R \phi \theta' \cot \theta +$$

$$\frac{1}{\sin^2 \theta} R \theta \phi'' = 0$$

÷ by $R \theta \phi$

$$\Rightarrow r^2 \frac{R''}{R} + 2r \frac{R'}{R} + \frac{\theta''}{\theta} + \frac{\theta'}{\theta} \cot \theta + \frac{\phi''}{\phi \sin^2 \theta} = 0$$

$$\Rightarrow r^2 \frac{R''}{R} + 2r \frac{R'}{R} = - \left[\frac{\theta''}{\theta} + \frac{\theta'}{\theta} \cot \theta + \frac{\phi''}{\phi \sin^2 \theta} \right] = n(n+1) \quad (\text{say})$$

Now,

$$r^2 \frac{R''}{R} + 2r \frac{R'}{R} = n(n+1) \quad \text{--- (3)}$$

$$\frac{\theta''}{\theta} + \frac{\theta'}{\theta} \cot \theta + \frac{1}{\sin^2 \theta} \frac{\phi''}{\phi} = -n(n+1) \quad \text{--- (4)}$$

$$\text{(3)} \Rightarrow r^2 R'' + 2r R' - n(n+1)R = 0$$

Consider, $z = \log r \Rightarrow e^z = r$

$$\text{Let } r^2 \frac{d^2}{dr^2} = D(D-1), \quad r \frac{d}{dr} = D, \quad e^z = r$$

$$\text{(3)} \Rightarrow [D(D-1) + 2D - n(n+1)]R = 0$$

$$[D^2 + D - n(n+1)]R = 0$$

Auxiliary equation

$$m^2 + m - n(n+1) = 0$$

$$m^2 + mn + m - n^2 - mn - n = 0$$

$$\Rightarrow m(m+n+1) - n(n+m+1) = 0$$

$$(m-n)(m+n+1) = 0$$

$$m = n, m = -(n+1)$$

$$\therefore R(z) = c_1 z^n + c_2 z^{-(n+1)}$$

$$R(r) = c_1 e^{n \log r} + c_2 e^{-(n+1) \log r}$$

$$= c_1 r^n + c_2 r^{-(n+1)}$$

$$R(r) = c_1 r^n + c_2 \frac{1}{r^{n+1}}$$

$$\textcircled{4} \Rightarrow \frac{\theta''}{\theta} + \cot \theta \frac{\theta'}{\theta} + \frac{1}{\sin^2 \theta} \frac{\phi''}{\phi} + n(n+1) = 0$$

ly by $\sin^2 \theta$

$$\sin^2 \theta \left\{ \frac{\theta''}{\theta} + \cot \theta \frac{\theta'}{\theta} + \frac{1}{\sin^2 \theta} \frac{\phi''}{\phi} + n(n+1) \right\} = 0 \textcircled{4}$$

and

$$\sin^2 \theta \left\{ \frac{\theta''}{\theta} + \cot \theta \frac{\theta'}{\theta} + n(n+1) \right\} = -\frac{\phi''}{\phi} = m^2$$

Now, we get,

$$\sin^2 \theta \left\{ \frac{\theta''}{\theta} + \cot \theta \frac{\theta'}{\theta} + n(n+1) \right\} = m^2 \textcircled{5} \quad \neq$$

$$\phi'' = -m^2 \phi \Rightarrow \phi'' + m^2 \phi = 0 \textcircled{6}$$

$$\textcircled{5} \Rightarrow \frac{\theta''}{\theta} + \cot \theta \frac{\theta'}{\theta} + n(n+1) = \frac{m^2}{\sin^2 \theta}$$

$$\Rightarrow \theta'' + \cot \theta \theta' + \left\{ n(n+1) - \frac{m^2}{\sin^2 \theta} \right\} \theta = 0$$

which is Legendre's equation

$$\therefore \Theta(\theta) = C_3 P_n^m \cos \theta + C_4 Q_n^m \cos \theta \quad \text{; result (i)}$$

where P_n^m and Q_n^m Legendre functions
1st and 2nd kind respectively and $Q_n \rightarrow \infty$ as $\theta \rightarrow \pi$

$$(b) \rightarrow (D^2 + m^2)\phi = 0$$

Auxiliary equation

$$P^2 + m^2 = 0$$

$$\therefore P = \pm im$$

$$\phi(\phi) = C_5 \cos m\phi + C_6 \sin m\phi$$

Here if $m=0$, then $\phi(\phi) = C_5$

i.e.) The equation (2) the solution of equation (1) does not depend on bounded variable.

This case is called axis symmetric case equation (2) becomes,

$$u(r, \theta, \phi) = \left(C_1 r^n + \frac{C_2}{r^{n+1}} \right) \cdot (C_3 P_n^m \cos \theta + C_4 Q_n^m \cos \theta) = (C_5 \cos m\phi + C_6 \sin m\phi)$$

For eliminating the unbounded solution choose $C_4 = 0$

$$u(r, \theta, \phi) = \left(C_1 r^n + \frac{C_2}{r^{n+1}} \right) (C_3 P_n^m \cos \theta)$$

$$(A \cos m\phi + B \sin m\phi)$$

In the axis symmetric case we have $m=0$.

Hence the above solution becomes

$$u(r, \theta, \phi) = (P_n^m \cos \theta) (A' r^n + B' / r^{n+1})$$

Thus the general solution of eqn (1) is given by

$$u(r, \theta, \phi) = \sum_{n=0}^{\infty} P_n^m \cos \theta \left(A_n r^n + \frac{B_n}{r^{n+1}} \right)$$

\therefore This is the general solution of spherical Co-Ordinates.

Interior Dirichlet Problem for circle

Statement:-

A Dirichlet problem for a circle is defined as follows.

To find the values of u at any point in the interior of the circle

$$\underline{r=a}$$

In terms of its values and the boundary $r=a$ such that,

u is a single valued and continuous within and on the circle region and satisfies the equation

$$\underline{\nabla^2 u = 0; \theta \leq r \leq a; 0 \leq \theta \leq 2\pi}$$

subject to $u(a, \theta) = f(\theta), 0 \leq \theta \leq 2\pi$.

Soln:-

The Laplace equation polar Co-Ordinate

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = 0 \quad \text{--- (1)}$$

Let $u(r, \theta) = R(r) \Theta(\theta)$ \rightarrow (2) be a solution of (1) sub (2) in (1)

$$\frac{d^2}{dr^2} [R(r)\Theta(\theta)] + \frac{1}{r} \frac{d}{dr} [R(r)\Theta(\theta)] + \frac{1}{r^2} \frac{d^2}{d\theta^2} [R(r)\Theta(\theta)] = 0$$

$$\Theta R'' + \frac{1}{r} R' \Theta + \frac{1}{r^2} R \Theta'' = 0$$

÷ by $R\Theta$.

$$\frac{R''}{R} + \frac{1}{r} \frac{R'}{R} + \frac{1}{r^2} \frac{\Theta''}{\Theta} = 0$$

× by r^2

$$r^2 \frac{R''}{R} + r \frac{R'}{R} + \frac{\Theta''}{\Theta} = 0$$

$$r^2 \frac{R''}{R} + r \frac{R'}{R} = -\frac{\Theta''}{\Theta} = k \text{ (say)}$$

$$\Rightarrow r^2 \frac{R''}{R} + r \frac{R'}{R} = k \quad \text{--- (3)}, \quad \Theta'' + k\Theta = 0 \quad \text{--- (4)}$$

$$\text{(3)} \Rightarrow r^2 R'' + rR' - kR = 0 \quad \text{--- (5)}$$

$$\text{(4)} \Rightarrow \Theta'' + k\Theta = 0 \quad \text{--- (6)}$$

Case (i): -

If $k > 0$, put $k = \lambda^2$

$$\text{(5)} \Rightarrow r^2 R'' + rR' - \lambda^2 R = 0$$

Let $z = \log r$; $r^2 \frac{d^2}{dr^2} = D(D-1)$, $r \frac{d}{dr} = D$

$$[D(D-1) + D - \lambda^2]R = 0$$

$$[D^2 - \lambda^2]R = 0$$

Auxiliary equation,

$$m^2 - \lambda^2 = 0$$

$$m = \pm \lambda$$

$$\therefore R(z) = C_1 e^{\lambda z} + C_2 e^{-\lambda z}$$

$$R(r) = c_1 e^{\lambda \log r} + c_2 e^{-\lambda \log r}$$

$$k(r) = c_1 r^\lambda + c_2 r^{-\lambda}$$

$$(6) \Rightarrow \frac{\theta''}{\theta} = -\lambda^2 \Rightarrow \theta'' + \lambda^2 \theta = 0$$

Auxillary equation:-

$$m^2 + \lambda^2 = 0$$

$$m = \pm i\lambda$$

$$\therefore \theta(\theta) = c_3 \cos \lambda \theta + c_4 \sin \lambda \theta$$

$$u(r, \theta) = (c_1 r^\lambda + c_2 r^{-\lambda}) \cdot (c_3 \cos \lambda \theta + c_4 \sin \lambda \theta) \quad (7)$$

Case (ii):-

If $k < 0$ put $k = -\lambda^2$

$$(5) \Rightarrow r^2 R'' + rR' + \lambda^2 R = 0$$

$$\Rightarrow [D(D-1) + D + \lambda^2]R = 0$$

$$[D^2 + \lambda^2]R = 0$$

Auxillary equation,

$$m^2 + \lambda^2 = 0$$

$$m = \pm i\lambda$$

$$R(z) = c_5 \cos \lambda z + c_6 \sin \lambda z$$

$$R(r) = c_5 \cos(\lambda \log r) + c_6 \sin(\lambda \log r)$$

$$(6) \Rightarrow \frac{\theta''}{\theta} = \lambda^2 \Rightarrow \theta'' - \lambda^2 \theta = 0$$

$$(D^2 - \lambda^2)\theta = 0$$

Auxillary equation

$$m^2 - \lambda^2 = 0$$

$$m = \pm \lambda$$

$$\therefore \theta(\theta) = c_7 e^{\lambda \theta} + c_8 e^{-\lambda \theta}$$

$$u(r, \theta) = [c_5 \cos(\lambda \log r) + c_6 \sin(\lambda \log r)]$$

$$(c_7 e^{\lambda \theta} + c_8 e^{-\lambda \theta}) \quad (8)$$

Then the required solution is obtained by the eqn (7)

$$i.e.] u(r, \theta) = (c_1 r^\lambda + c_2 r^{-\lambda}) (c_3 \cos \lambda \theta + c_4 \sin \lambda \theta)$$

Let $r=0$, $u(r, \theta)$ must be finite and hence we can choose $c_2 = 0$.

$$u(r, \theta) = (c_1 r^\lambda) (c_3 \cos \lambda \theta + c_4 \sin \lambda \theta).$$

$$u(r, \theta) = (A \cos \lambda \theta + B \sin \lambda \theta) r^\lambda.$$

We know that,

The Periodically Condition

$$u(r, \theta) = u(r, \theta + 2\pi)$$

$$u(r, \theta) = r^\lambda [A \cos \lambda (\theta + 2\pi) + B \sin \lambda (\theta + 2\pi)]$$

$$= r^\lambda [A \cos (\lambda \theta + 2\pi \lambda) + B \sin (\lambda \theta + 2\pi \lambda)]$$

$$= r^\lambda [A (\cos \lambda \theta \cdot \cos 2\pi \lambda - \sin \lambda \theta \cdot \sin 2\pi \lambda)$$

$$+ B (\sin \lambda \theta \cdot \cos 2\pi \lambda + \cos \lambda \theta \cdot \sin 2\pi \lambda)]$$

$$(A \cos \lambda \theta + B \sin \lambda \theta) r^\lambda = r^\lambda [A (\cos \lambda \theta \cdot \cos 2\pi \lambda - \sin \lambda \theta \cdot \sin 2\pi \lambda)$$

$$+ B (\sin \lambda \theta \cdot \cos 2\pi \lambda + \cos \lambda \theta \cdot \sin 2\pi \lambda)]$$

$$\Rightarrow A [\cos \lambda \theta - \cos \lambda \theta \cdot \cos 2\pi \lambda + \sin \lambda \theta \cdot \sin 2\pi \lambda]$$

$$+ B (\sin \lambda \theta - \sin \lambda \theta \cdot \cos 2\pi \lambda - \cos \lambda \theta \cdot \sin 2\pi \lambda) = 0$$

$$\sin^2 2\theta = 2 \sin \theta \cos \theta$$

$$1 - \cos^2 \theta = 2 \sin^2 \theta$$

$$\Rightarrow A [\cos \lambda \theta (1 - \cos 2\pi \lambda) + \sin \lambda \theta \cdot \sin 2\pi \lambda] +$$

$$B [\sin \lambda \theta (1 - \cos 2\pi \lambda) - \cos \lambda \theta \cdot \sin 2\pi \lambda] = 0$$

$$\Rightarrow A [2 \cos \lambda \theta \cdot \sin^2 \lambda \pi + 2 \sin \lambda \theta \cdot \sin \pi \lambda \cdot \cos \pi \lambda]$$

$$+ B [2 \sin \lambda \theta \cdot \sin^2 \lambda \pi - 2 \sin \pi \lambda \cdot \cos \pi \lambda \cdot \cos \lambda \theta] = 0$$

$$\Rightarrow 2 \sin \lambda \pi \left[A (\cos \lambda \theta \cdot \sin \lambda \pi + \sin \lambda \theta \cdot \cos \lambda \pi) + B (\sin \lambda \theta \cdot \sin \lambda \pi - \cos \lambda \theta \cdot \cos \lambda \pi) \right] = 0$$

$$\Rightarrow 2 \sin \lambda \pi \left\{ A \sin (\lambda \theta + \lambda \pi) + B (-\cos (\lambda \theta + \lambda \pi)) \right\} = 0$$

$$\Rightarrow 2 \sin \lambda \pi \left\{ A \sin (\lambda \theta + \lambda \pi) + B (-\cos (\lambda \theta + \lambda \pi)) \right\} = 0$$

$$2 \sin \lambda \pi = 0$$

$$\lambda \pi = n \pi ; n = 0, 1, 2, \dots$$

$$\lambda = n$$

$$\therefore u(r, \theta) = (A \cos n \theta + B \sin n \theta) r^n ; n = 0, 1, 2, \dots$$

In general,

$$u(r, \theta) = \sum_{n=0}^{\infty} r^n (A_n \cos n \theta + B_n \sin n \theta)$$

$$u(r, \theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} r^n (a_n \cos n \theta + b_n \sin n \theta) \quad \text{--- (10)}$$

Where $A_0 = a_0/2$; $A_n = a_n$; $B_n = b_n$ $\forall n = 1, 2, \dots$

Which is the full range fourier series,

Given that

$$u(a, \theta) = f(\theta)$$

$$u(a, \theta) = f(\theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} (r^n a_n \cos n \theta + r^n b_n \sin n \theta)$$

$$u(a, \theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos n \theta + b_n \sin n \theta) \quad \text{--- (11)}$$

Where the fourier co-efficient series

are

$$a_0 = \frac{1}{\pi} \int_0^{2\pi} f(\phi) d\phi ; a_n = \frac{1}{\pi} \int_0^{2\pi} f(\phi) \cos^n \phi d\phi$$

$$a_n = \frac{1}{a^n \pi} \int_0^{2\pi} f(\phi) \cos^n \phi d\phi$$

$$b_n = \frac{1}{a^n \pi} \int_0^{2\pi} f(\phi) \sin^n \phi d\phi$$

$$x + x^2 + x^3 + \dots \\ x(1 + x + x^2 + \dots)$$

From eqn (10)

$$u(r, \theta) = \frac{1}{2\pi} \int_0^{2\pi} f(\phi) d\phi + \sum_{n=1}^{\infty} r^n \left\{ \frac{\cos n\theta}{\pi a^n} \int_0^{2\pi} f(\phi) \cos n\phi d\phi + \frac{\sin n\theta}{\pi a^n} \int_0^{2\pi} f(\phi) \sin n\phi d\phi \right\}$$

$$= \frac{1}{\pi} \left[\int_0^{2\pi} \left(\frac{f(\phi)}{2} + \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n f(\phi) \cos n\phi \cos n\theta + \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n f(\phi) \sin n\phi \sin n\theta \right) d\phi \right]$$

$$u(r, \theta) = \frac{1}{\pi} \int_0^{2\pi} f(\phi) \left[\frac{1}{2} + \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \cos n(\phi - \theta) \right] d\phi$$

$$\left[\cos n\phi \cdot \frac{\cos(A-B)}{\cos n\theta + \sin n\phi \sin n\theta} \right] d\phi$$

$$u(r, \theta) = \frac{1}{\pi} \int_0^{2\pi} f(\phi) \left\{ \frac{1}{2} + \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \cos n(\phi - \theta) \right\} d\phi \quad \text{--- (12)}$$

$$\text{Let } C = \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \cos n(\phi - \theta)$$

$$\text{Assume, } S = \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \sin n(\phi - \theta)$$

$$C + iS = \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \left\{ \cos n(\phi - \theta) + i \sin n(\phi - \theta) \right\}$$

$$= \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n e^{in(\phi - \theta)}$$

$$e^{i\theta} = \cos \theta + i \sin \theta$$

$$= \sum_{n=1}^{\infty} \left[\left(\frac{r}{a} \right) e^{i(\phi - \theta)} \right]^n$$

$$= \frac{r}{a} e^{i(\phi-\theta)} + \left[\left(\frac{r}{a}\right) e^{i(\phi-\theta)} \right]^2 + \dots$$

$$= \left(\frac{r}{a} e^{i(\phi-\theta)}\right) \left[1 - \left(\frac{r}{a} e^{i(\phi-\theta)}\right) \right]^{-1}$$

$$S^{\infty} = \left(\frac{r}{a}\right) e^{i(\phi-\theta)} \left[1 - \left(\frac{r}{a}\right) e^{i(\phi-\theta)} \right]^{-1}$$

$$C + iS = \frac{\left(\frac{r}{a}\right) e^{i(\phi-\theta)}}{\left(1 - \left(\frac{r}{a}\right) e^{i(\phi-\theta)}\right)} \times \frac{1 - \left(\frac{r}{a}\right) e^{-i(\phi-\theta)}}{1 - \left(\frac{r}{a}\right) e^{-i(\phi-\theta)}}$$

$$= \frac{\left(\frac{r}{a}\right) e^{i(\phi-\theta)} - \left(\frac{r}{a}\right)^2}{\dots}$$

$$\left[1 - \left(\frac{r}{a}\right) (\cos(\phi-\theta) + i \sin(\phi-\theta)) \right]$$

$$\left[1 - \frac{r}{a} (\cos(\phi-\theta) - i \sin(\phi-\theta)) \right]$$

$$= \frac{\left(\frac{r}{a}\right) [\cos(\phi-\theta) + i \sin(\phi-\theta)] - \left(\frac{r}{a}\right)^2}{\dots}$$

$$\left[1 - \left(\frac{r}{a}\right) \cos(\phi-\theta) \right]^2 + \left[\left(\frac{r}{a}\right) \sin(\phi-\theta) \right]^2$$

$$= \frac{\left(\frac{r}{a}\right) [\cos(\phi-\theta) + i \sin(\phi-\theta)] - \left(\frac{r}{a}\right)^2}{\dots}$$

$$1 + \frac{r^2}{a^2} \cos^2(\phi-\theta) - \frac{2r}{a} \cos(\phi-\theta) + \frac{r^2}{a^2} \sin^2(\phi-\theta)$$

$$C + iS = \frac{\left(\frac{r}{a}\right) \cos(\phi-\theta) - \frac{r^2}{a^2} + i \left(\frac{r}{a}\right) \sin(\phi-\theta)}{\dots}$$

$$1 + \frac{r^2}{a^2} - 2\left(\frac{r}{a}\right) \cos(\phi-\theta)$$

Equating real part,

$$c = \frac{\left(\frac{r}{a}\right) \cos(\phi - \theta) - \frac{r^2}{a^2}}{1 + \frac{r^2}{a^2} - 2\frac{r}{a} \cos(\phi - \theta)}$$

$$c + \frac{1}{2} = \frac{\left(\frac{r}{a}\right) \cos(\phi - \theta) - \frac{r^2}{a^2}}{1 + \frac{r^2}{a^2} - 2\frac{r}{a} \cos(\phi - \theta)} + \frac{1}{2}$$

$$c + \frac{1}{2} = \frac{2 \left[\cancel{\left(\frac{r}{a}\right) \cos(\phi - \theta) - \frac{r^2}{a^2}} \right] + \left[1 + \frac{r^2}{a^2} - \frac{2r}{a} \cos(\phi - \theta) \right]}{2 \left[1 + \frac{r^2}{a^2} - 2\frac{r}{a} \cos(\phi - \theta) \right]}$$

$$= \frac{1 - \left(\frac{r}{a}\right)^2}{2 \left[1 + \left(\frac{r}{a}\right)^2 - 2\frac{r}{a} \cos(\phi - \theta) \right]}$$

$$= \frac{1 - \left(\frac{r}{a}\right)^2}{\frac{2}{a^2} (a^2 + r^2 - 2ar \cos(\phi - \theta))}$$

$$= \frac{a^2 - r^2}{2(a^2 + r^2 - 2ar \cos(\phi - \theta))} \quad \text{--- (13)}$$

Sub (13) on (2)

$$u(r, \theta) = \frac{1}{2\pi} \int_0^{2\pi} f(\phi) \left[\frac{a^2 - r^2}{a^2 - 2ar \cos(\phi - \theta) + r^2} \right] d\phi$$

This is called a interior dirichlet

Problem for a circle.

✓ Exterior Dirichlet Problem for Circle:-

statement:-
57)

The Exterior Dirichlet problem is described by the $\nabla^2 \phi = 0$ for $0 \leq \theta \leq 2\pi$, $r \geq a$ and $\phi(a, \theta) = f(\theta)$ where $r = a$. where $f(\theta)$ is continuous function of θ at $r = a$ and ϕ must be bounded as $r \rightarrow \infty$.

Proof:-

$$\phi(r, \theta) = \sum_{n=0}^{\infty} (C_n r^n + D_n r^{-n}) (A_n \cos n\theta + B_n \sin n\theta)$$

Now, as $r \rightarrow \infty$ we require ϕ to be bounded. \therefore Take $C_n = 0$

$$\phi(r, \theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} r^{-n} (a_n \cos n\theta + b_n \sin n\theta) \quad D_n = 1$$

Now $a_0 = \frac{1}{2\pi} \int_0^{2\pi} f(\theta) d\theta$

$$\phi(r, \theta) = \frac{a}{r} e^{i(\phi-\theta)} + \left[\left(\frac{a}{r}\right) e^{i(\phi-\theta)} \right]^2 + \dots$$

$$= \left(\frac{a}{r}\right) e^{i(\phi-\theta)} \left[1 + \left(\frac{a}{r}\right) e^{i(\phi-\theta)} + \dots \right]$$

$$= \left(\frac{a}{r}\right) e^{i(\phi-\theta)} \left[1 - \left(\frac{a}{r}\right) e^{i(\phi-\theta)} \right]^{-1}$$

$$= \frac{\left(\frac{a}{r}\right) e^{i(\phi-\theta)}}{1 - \left(\frac{a}{r}\right) e^{i(\phi-\theta)}}$$

$$= \frac{\left(\frac{a}{r}\right) e^{i(\phi-\theta)}}{1 - \left(\frac{a}{r}\right) e^{i(\phi-\theta)}} \times \frac{1 - \left(\frac{a}{r}\right) e^{-i(\phi-\theta)}}{1 - \left(\frac{a}{r}\right) e^{-i(\phi-\theta)}}$$

by (2) $\phi(r, \theta) = \frac{1}{2\pi} \int_0^{2\pi} f(\theta) \left[\frac{1}{2} + \sum_{n=1}^{\infty} \left(\frac{a}{r}\right)^n \cos n(\theta - \phi) \right] d\theta$

$$= \frac{\left(\frac{a}{r}\right) e^{i(\phi-\theta)} - \left(a/r\right)^2}{\left(1 - \left(\frac{a}{r}\right) e^{i(\phi-\theta)}\right) \left[1 - \left(\frac{a}{r}\right) e^{-i(\phi-\theta)}\right]}$$

$$\phi(r, \theta) = \frac{\left(\frac{a}{r}\right) e^{i(\phi-\theta)} - \left(a/r\right)^2}{\left[1 - \left(\frac{a}{r}\right) \left[\cos(\phi-\theta) + i \sin(\phi-\theta)\right]\right] \left[1 - \left(\frac{a}{r}\right) \left[\cos(\phi-\theta) - i \sin(\phi-\theta)\right]\right]}$$

$$= \frac{\left(\frac{a}{r}\right) [\cos(\phi-\theta) + i \sin(\phi-\theta)] - \left(a/r\right)^2}{\left[1 - \left(\frac{a}{r}\right) \cos(\phi-\theta)\right]^2 + \left[\left(\frac{a}{r}\right) \sin(\phi-\theta)\right]^2}$$

(a-b)² formula.

$$= \frac{\left(\frac{a}{r}\right) \cos(\phi-\theta) - \left(\frac{a}{r}\right)^2 + i \left(\frac{a}{r}\right) \sin(\phi-\theta)}{1 + \left(\frac{a}{r}\right)^2 \cos^2(\phi-\theta) - 2\left(\frac{a}{r}\right) \cos(\phi-\theta) + \frac{a^2}{r^2} \sin^2(\phi-\theta)}$$

$$= \frac{\left(\frac{a}{r}\right) \cos(\phi-\theta) - \left(a/r\right)^2 + i \left(\frac{a}{r}\right) \sin(\phi-\theta)}{1 + \left(\frac{a}{r}\right)^2 - 2\left(\frac{a}{r}\right) \cos(\phi-\theta)}$$

Equating the real parts

$$c = \frac{\left(\frac{a}{r}\right) \cos(\phi-\theta) - \left(\frac{a}{r}\right)^2}{1 + \left(\frac{a}{r}\right)^2 - 2\left(\frac{a}{r}\right) \cos(\phi-\theta)}$$

$$c + \frac{1}{2} = \frac{\left(\frac{a}{r}\right) \cos(\phi-\theta) - \left(\frac{a}{r}\right)^2}{1 + \left(\frac{a}{r}\right)^2 - 2\left(\frac{a}{r}\right) \cos(\phi-\theta)} + \frac{1}{2}$$

$$= \frac{2\left(\frac{a}{r}\right) \cos(\phi-\theta) - 2\frac{a^2}{r^2} + 1 + \frac{a^2}{r^2} - 2\left(\frac{a}{r}\right) \cos(\phi-\theta)}{2 \left[1 + \left(\frac{a}{r}\right)^2 - 2\left(\frac{a}{r}\right) \cos(\phi-\theta)\right]}$$

$$c + \frac{1}{2} = \frac{1 - a^2/r^2}{2 \left[1 + \left(\frac{a}{r}\right)^2 - 2\left(\frac{a}{r}\right) \cos(\phi-\theta)\right]}$$

$$= \frac{r^2 - a^2}{2[r^2 + a^2 - 2ar \cos(\phi - \theta)]}$$

Sub (6) in (5)

$$\phi(r, \theta) = \frac{1}{2\pi} \int_0^{2\pi} f(\phi) \left[\frac{r^2 - a^2}{r^2 + a^2 - 2ar \cos(\phi - \theta)} \right] d\phi$$

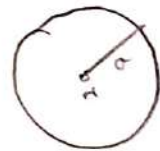
This is the solution of exterior Dirichlet for a circle.

Interior Neumann for a Circle: -

Statement:-

Q2M (Find the value of u at any point in the interior of the circle $r=a$ such that $\nabla^2 u = 0$, $0 \leq r \leq a$ and $\frac{\partial u}{\partial n} - \frac{\partial u}{\partial r} = g(\theta)$ on $r=a$.)

Proof:- Given $\nabla^2 u = 0$ — (1)



$$\frac{\partial u}{\partial n} - \frac{\partial u}{\partial r} = g(\theta) \text{ — (2)}$$

Then by the method of separation of variables the general solution of equation (1) is given by

$$u(r, \theta) = \sum_{n=0}^{\infty} (C_n r^n + D_n r^{-n}) (A_n \cos n\theta + B_n \sin n\theta)$$

This gives u is unbounded where $r=0$

where $D_n = 0$

$$\begin{aligned} u(r, \theta) &= \sum_{n=0}^{\infty} C_n r^n (A_n \cos n\theta + B_n \sin n\theta) \\ &= \sum_{n=0}^{\infty} r^n (A_n \cos n\theta + B_n \sin n\theta) \end{aligned}$$

$$= A_0 + \sum_{n=1}^{\infty} r^n (A_n \cos n\theta + B_n \sin n\theta)$$

$$u(r, \theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} r^n (a_n \cos n\theta + b_n \sin n\theta) \quad \text{--- (3)}$$

where $A_0 = \frac{a_0}{2}$, $A_n = a_n$, $B_n = b_n$ for $n=1, 2, \dots$
 differentiable partially w.r. to r .

$$\frac{\partial u}{\partial r} = 0 + \sum_{n=1}^{\infty} n r^{n-1} (a_n \cos n\theta + b_n \sin n\theta)$$

Given that $\frac{\partial u(a, \theta)}{\partial r} = g(\theta)$, where $r=a$

$$\Rightarrow g(\theta) = \left(\frac{\partial u}{\partial r} \right)_{r=a}$$

$$= \sum_{n=1}^{\infty} n a^{n-1} (a_n \cos n\theta + b_n \sin n\theta)$$

$$g(\theta) = \sum_{n=1}^{\infty} (n a^{n-1} a_n \cos n\theta + b_n \sin n\theta)$$

which is the full range Fourier series.
 whose Fourier coefficients are, $a_0 = 0$.

$$n a^{n-1} a_n = \frac{1}{\pi} \int_0^{2\pi} g(\phi) \cos n\phi d\phi$$

$$\Rightarrow a_n = \frac{1}{n a^{n-1} \pi} \int_0^{2\pi} g(\phi) \cos n\phi d\phi \quad \text{--- (i)}$$

$$n a^{n-1} b_n = \frac{1}{\pi} \int_0^{2\pi} g(\phi) \sin n\phi d\phi$$

$$b_n = \frac{1}{n a^{n-1} \pi} \int_0^{2\pi} g(\phi) \sin n\phi d\phi \quad \text{--- (ii)}$$

Sub (i) and (ii) in (3)

$$u(r, \theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} r^n \left[\frac{1}{n a^{n-1} \pi} \int_0^{2\pi} g(\phi) \cos n\phi \cdot \cos n\theta d\phi + \frac{1}{n a^{n-1} \pi} \int_0^{2\pi} g(\phi) \sin n\phi \cdot \sin n\theta d\phi \right]$$

$$\phi(r, \theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} r^n \frac{1}{n a^{n-1} \pi} \left[\int_0^{2\pi} g(\phi) (\cos n\phi \cos n\theta + \sin n\phi \sin n\theta) d\phi \right]$$

$$= \frac{a_0}{2} + \sum_{n=1}^{\infty} r^n \frac{1}{n a^{n-1} \pi} \left[\int_0^{2\pi} g(\phi) \cos n(\phi - \theta) d\phi \right]$$

$$= \frac{a_0}{2} + \sum_{n=1}^{\infty} r^n \frac{a}{n a^n \pi} \left[\int_0^{2\pi} g(\phi) \cdot \cos n(\phi - \theta) d\phi \right]$$

$$\phi(r, \theta) = \frac{a_0}{2} + \frac{a}{\pi} \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \frac{1}{n} \int_0^{2\pi} g(\phi) \cos n(\phi - \theta) d\phi$$

$$= \frac{a_0}{2} + \frac{a}{\pi} \int_0^{2\pi} g(\phi) \left[\sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \left(\frac{1}{n} \right) \cos n(\phi - \theta) \right] d\phi \quad \text{--- (A)}$$

$$\text{Let } C = \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \left(\frac{1}{n} \right) \cos n(\phi - \theta)$$

$$S = \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \left(\frac{1}{n} \right) \sin n(\phi - \theta)$$

$$\left[\frac{\partial u}{\partial n} - \frac{\partial u}{\partial r} = g(\theta) \text{ on } r=a \right]$$

$$C + iS = \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \left(\frac{1}{n} \right) [\cos n(\phi - \theta) + i \sin n(\phi - \theta)]$$

$$= \sum_{n=1}^{\infty} \left(\frac{r}{a} \right)^n \left(\frac{1}{n} \right) e^{in(\phi - \theta)}$$

$e^{in\theta}$

$$= \sum_{n=1}^{\infty} \left[\left(\frac{r}{a} \right) e^{i(\phi - \theta)} \right]^n \frac{1}{n}$$

$$= \left(\frac{r}{a} \right) e^{i(\phi - \theta)} + \left[\left(\frac{r}{a} \right) e^{i(\phi - \theta)} \right]^2 \frac{1}{2} + \left[\left(\frac{r}{a} \right) e^{i(\phi - \theta)} \right]^3 \frac{1}{3} + \dots$$

$$c + is = -\log \left[1 - \left(\frac{r}{a} \right) e^{i(\phi - \theta)} \right]$$

where $z = 1 - \left(\frac{r}{a} \right) e^{i(\phi - \theta)}$

$$e^{c+is} = e^{-\log z}$$

$$= z^{-1}$$

$$= \frac{1}{z}$$

$$z = \frac{1}{e^{c+is}} = e^{-(c+is)}$$

$$= e^{-c} \cdot e^{-is}$$

$$|z| = |e^{-c}| |e^{-is}|$$

Take real part

$$|z| = |e^{-c}| = e^{-c}$$

$$\Rightarrow \log |z| = \log e^{-c}$$

$$\Rightarrow \log |z| = -c$$

$$c = -\log |z| \Rightarrow c = -\log \left| 1 - \left(\frac{r}{a} \right) e^{i(\phi - \theta)} \right|$$

$$= -\log \left| 1 - \left(\frac{r}{a} \right) (\cos(\phi - \theta) + i \sin(\phi - \theta)) \right|$$

$$= -\log \left| \left(1 - \left(\frac{r}{a} \right) \cos(\phi - \theta) \right) - \left(\frac{ir}{a} \right) \sin(\phi - \theta) \right|$$

$$= -\log \left| \left(1 - \left(\frac{r}{a} \right) \cos(\phi - \theta) \right)^2 + \left(\frac{r}{a} \right)^2 \sin^2(\phi - \theta) \right|$$

$$= -\log \sqrt{\left(1 - \left(\frac{r}{a} \right) \cos(\phi - \theta) \right)^2 + \left(\frac{r}{a} \right)^2 \sin^2(\phi - \theta)}$$

$$= -\log \sqrt{1 - \frac{2r}{a} \cos(\phi - \theta) + \left(\frac{r}{a} \right)^2}$$

$$\log m^n = n \log m$$

$$\log(1-x) = -x - \frac{x^2}{2} - \frac{x^3}{3} - \dots$$

$$= -\log \left[\sqrt{\frac{1}{a^2} (a^2 - 2ar \cos(\phi - \theta) + r^2)} \right]$$

$$= -\log \left(\frac{1}{a} \right) \sqrt{a^2 + r^2 - 2ar \cos(\phi - \theta)} \quad \text{--- (E)}$$

Sub in eqn (A)

$$\phi(r, \theta) = \frac{a_0}{2} + \frac{a}{\pi} \int_0^{2\pi} g(\phi) \left[-\ln \left(\frac{1}{a} \right) \sqrt{a^2 + r^2 - 2ar \cos(\phi - \theta)} \right] d\phi$$

$$\phi(r, \theta) = \frac{a_0}{2} - \frac{a}{\pi} \int_0^{2\pi} g(\phi) \left[\ln \left(\frac{1}{a} \right) \sqrt{a^2 + r^2 - 2ar \cos(\phi - \theta)} \right] d\phi$$

This is the interior Neumann for a circle.

RESULT:-

$$\int_{-1}^1 [P_n(\mu)]^2 d\mu = \frac{2}{2n+1} \frac{(n!)^2}{(n-n)!}$$

Interior Dirichlet problem for a sphere:-

Statement:-

Find the value of u at any point in the interior of sphere $r=a$ such that

$$\nabla^2 u = 0 \rightarrow (1) \quad 0 \leq r \leq a, \quad 0 \leq \theta \leq 2\pi \text{ and}$$

$$u(a, \theta, \phi) = f(\theta, \phi) \text{ on } r=a.$$

Proof:-

By using separation variable method.

The general solution of eqn (1) in spherical co-ordinates is given by

$$u(r, \theta, \phi) = \left[c_1 r^n + \frac{c_2}{r^{n+1}} \right] \left[c_3 P_n^m \cos \theta + c_4 Q_n^m \cos \theta \right] \\ \left[c_5 \cos m\phi + c_6 \sin m\phi \right] \quad \text{--- (2)}$$

where $P_n^m \cos \theta$ and $Q_n^m \cos \theta$ are Legendre function of 1st and 2nd kind respectively.

Since $Q_n^m \cos \theta \rightarrow \infty$ as $\theta \rightarrow 0$.

We can choose $c_4 = 0$.

Since eqn (2) is unbounded when $r = 0$.

$$\Rightarrow u(r, \theta, \phi) = r^n P_n^m \cos \theta \left[A_{nm} \cos m\phi + B_{nm} \sin m\phi \right]$$

using given boundary condition

$$u(a, \theta, \phi) = f(\theta, \phi) \text{ where } r = a.$$

$$\Rightarrow f(\theta, \phi) = u(a, \theta, \phi)$$

$$f(\theta, \phi) = \sum_{n=0}^{\infty} \sum_{m=0}^n a^n P_n^m \cos \theta \left[A_{nm} \cos m\phi + B_{nm} \sin m\phi \right]$$

--- (3)

In Particular,

$$f(\theta, \phi) = a^n P_n^m \cos \theta \left[A_{nm} \cos m\phi + B_{nm} \sin m\phi \right]$$

Let $\mu = \cos \theta$.

$$f(\theta, \phi) = a^n P_n^m(\mu) \left[A_{nm} \cos m\phi + B_{nm} \sin m\phi \right]$$

Integrate both sides with respect to μ ,

$-1 \leq \mu \leq 1$ and multiply by $\cos m\phi P_n^m(\mu)$

We get,

$$\int_{-1}^1 f(\theta, \phi) \cos m\phi P_n^m(\mu) d\mu = \int_{-1}^1 a^n \left[P_n^m(\mu) \right]^2 \left[A_{nm} \cos^2 m\phi + B_{nm} \cos m\phi \cdot \sin m\phi \right] d\mu$$

$$2 a^n \left[A_{mn} \cos^2 m\phi + B_{mn} \sin m\phi \cdot \cos m\phi \right] \left[P_n^m(\mu) \right] d\mu$$

$\cos^2 = \frac{1 + \cos 2\phi}{2}$

$$\left[A_{mn} \left(\frac{1 + \cos 2m\phi}{2} \right) + B_{mn} \left(\frac{\sin 2m\phi}{2} \right) \right]$$

$$\left[\frac{2}{2n+1} \cdot \frac{(m+n)!}{(m-n)!} \right] \cdot a^n$$

Integrate with respect to ϕ , $0 \leq \phi \leq 2\pi$.

We have,

$$\int_0^{2\pi} \int_{-1}^1 f(\theta, \phi) \cos m\phi P_n^m(\mu) d\mu d\phi = \frac{2a^n (m+n)!}{(2n+1)(m-n)!} \int_0^{2\pi} \left[A_{mn} \left(\frac{1 + \cos 2m\phi}{2} \right) + B_{mn} \left(\frac{\sin 2m\phi}{2} \right) \right] d\phi$$

$$= \frac{2a^n (m+n)!}{(2n+1)(m-n)!} \left\{ \frac{A_{mn}}{2} \left[\phi + \frac{\sin 2m\phi}{2m} \right]_0^{2\pi} + \frac{B_{mn}}{2} \left[-\frac{\cos 2m\phi}{2m} \right]_0^{2\pi} \right\}$$

$\sin 2\pi = 0$

$$= \frac{2a^n (m+n)!}{(2n+1)(m-n)!} \left[\frac{A_{mn}}{2} (2\pi) + \frac{B_{mn}}{2} \left(-\frac{1}{2m} + \frac{1}{2m} \right) \right]$$

$$= \frac{2a^n (m+n)!}{(2n+1)(m-n)!} A_{mn} \pi$$

$$A_{mn} = \frac{1}{2\pi} \frac{(2n+1)(m-n)!}{a^n (m+n)!} \int_0^{2\pi} \int_{-1}^1 f(\theta, \phi) \cos m\phi P_n^m(\mu) d\mu d\phi$$

$$B_{mn} = \frac{1}{2\pi} \frac{(2n+1)(m-n)!}{a^n (m+n)!} \int_0^{2\pi} \int_{-1}^1 f(\theta, \phi) \sin m\phi P_n^m(\mu) d\mu d\phi$$

eqn (3) \Rightarrow

$$u(r, \theta, \phi) = \sum_{n=0}^{\infty} \sum_{m=0}^n \gamma_n^m P_n^m(\mu) \left[\frac{(2n+1)(m-n)!}{(2\pi a^n (m+n)!)} \int_0^{2\pi} \int_{-1}^1 f(\psi, \psi) \cos m\psi d\psi d\psi \right]$$

$$\cos m\phi P_n^m(\mu) d\phi d\psi + \frac{(2n+1)(m-n)!}{2\pi a^n (m+n)!} \int_0^{2\pi} \int_{-1}^1 f(\psi, \gamma) \sin m\psi P_n^m(\mu') \sin n\gamma d\psi d\gamma.$$

$$= \frac{1}{2\pi} \int_0^{2\pi} \int_{-1}^1 f(\psi, \gamma) \sum_{n=0}^{\infty} \sum_{m=0}^n \left(\frac{\gamma}{a}\right)^n \frac{(2n+1)(m-n)!}{(m+n)!} P_n^m(\mu')$$

$$P_n^m(\mu') [\cos m\psi \cos m\phi + \sin m\psi \sin m\phi] d\psi d\phi$$

$$= \frac{1}{2\pi} \int_0^{2\pi} \int_{-1}^1 f(\psi, \gamma) \sum_{n=0}^{\infty} \sum_{m=0}^n \left(\frac{\gamma}{a}\right)^n \frac{(2n+1)(m-n)!}{(m+n)!}$$

$$P_n^m \cos \theta \cdot P_n^m(\cos \psi) \cos m(\psi - \phi) d\psi d\gamma$$

$$u(r, \theta, \phi) = \frac{1}{2\pi} \sum_{n=0}^{\infty} \sum_{m=0}^n \left(\frac{\gamma}{a}\right)^n \frac{(2n+1)(m-n)!}{(m+n)!} \int_0^{2\pi} \int_{-1}^1$$

$$f(\psi, \gamma) P_n^m \cos \theta \cdot P_n^m(\cos \psi) \cos m(\psi - \phi) d\psi d\gamma.$$

which is the interior Dirichlet problem for a sphere.

